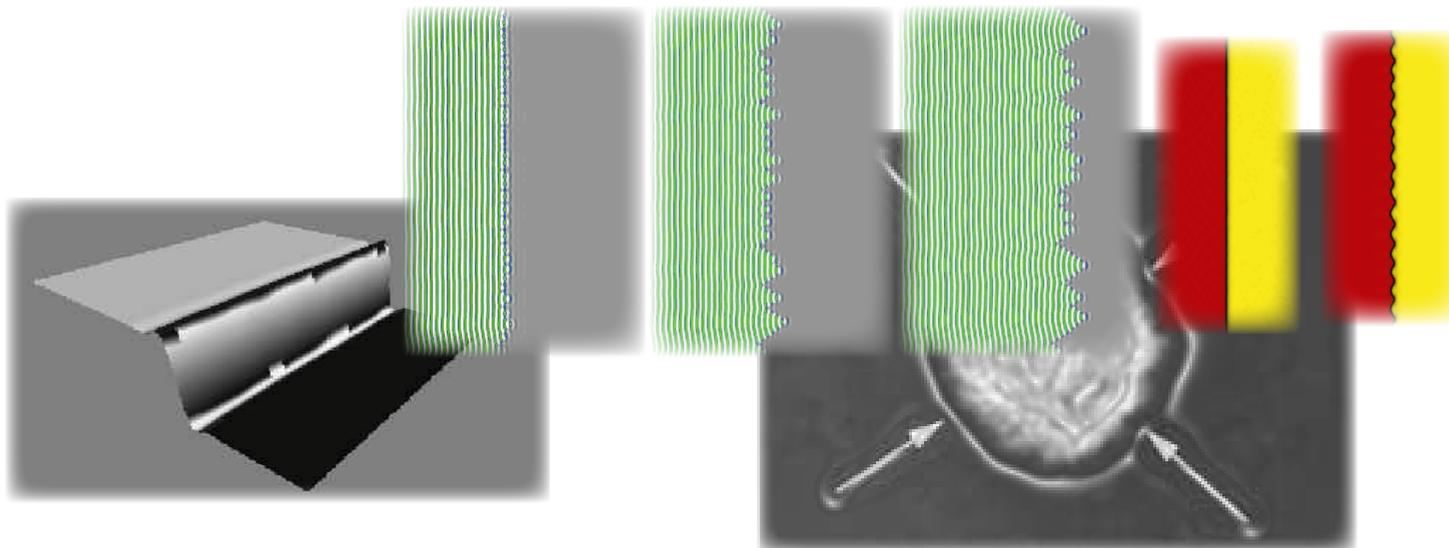


Interface dynamics of a front connecting stripe pattern with uniform state

dfi Departamento de Física.
Facultad de ciencias
Físicas y Matemáticas,
Universidad de Chile.

Marcel G. Clerc, Daniel Escaff and René Rojas

*Departamento de Física, FCFM,
Universidad de Chile.*



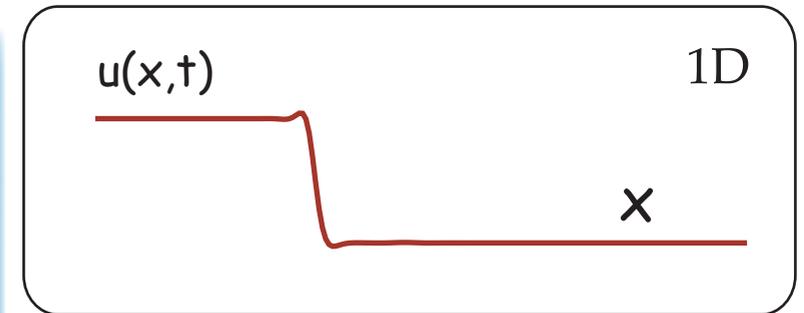
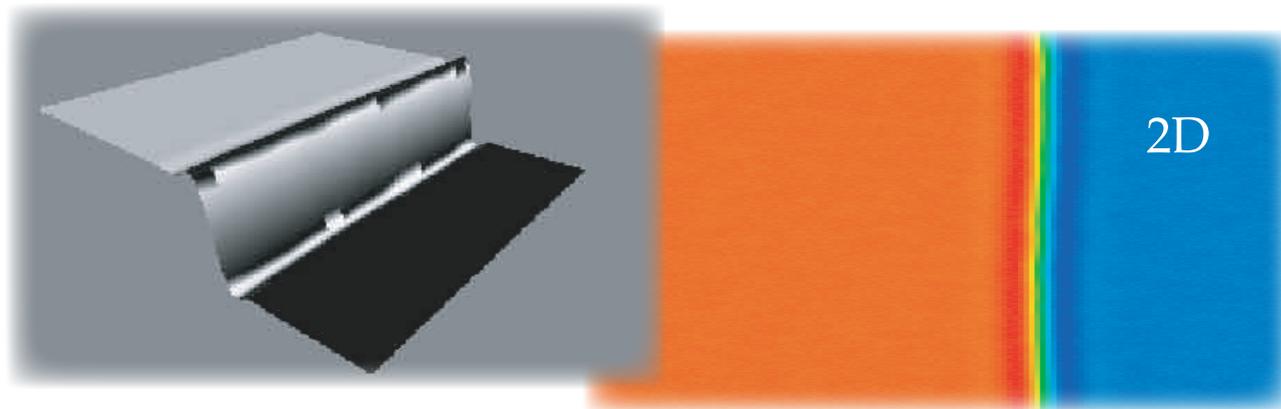
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Universidad de Chile.

Ring Program ACT 15, PBCT & CFMAT

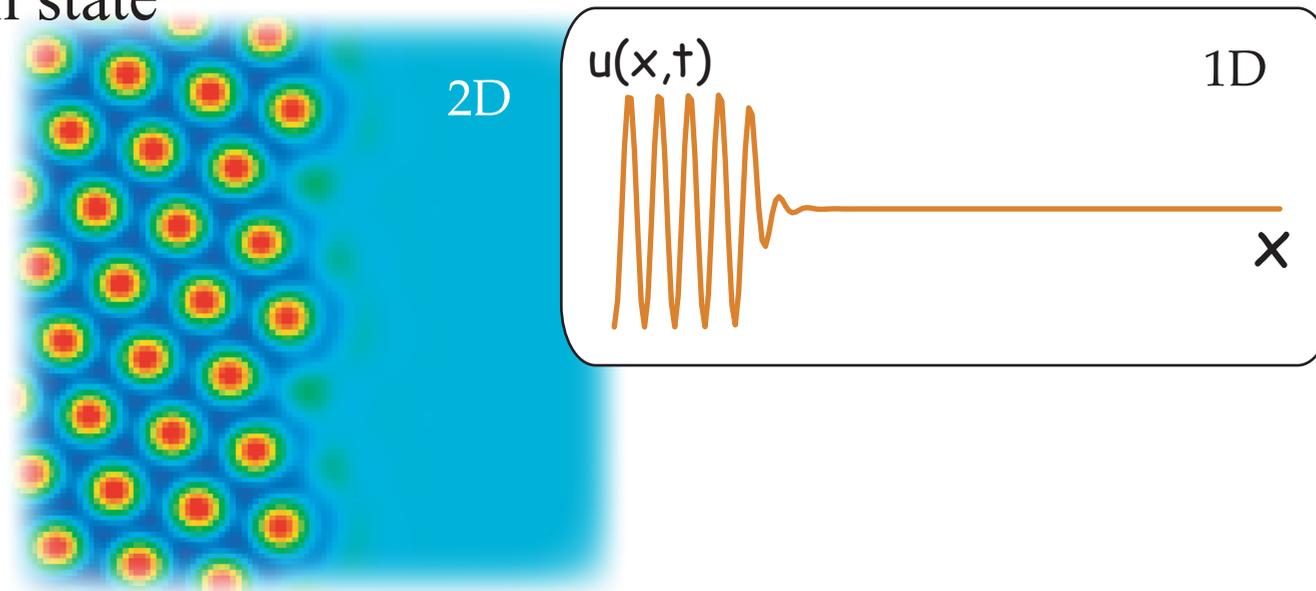
Front solution

- "Stationary solution that links two steady states".
- "In 1D the front solution are heroclinic connection of stationaries states in the stationary extended system or moving reference frame" (multi-stability).

Two uniform states

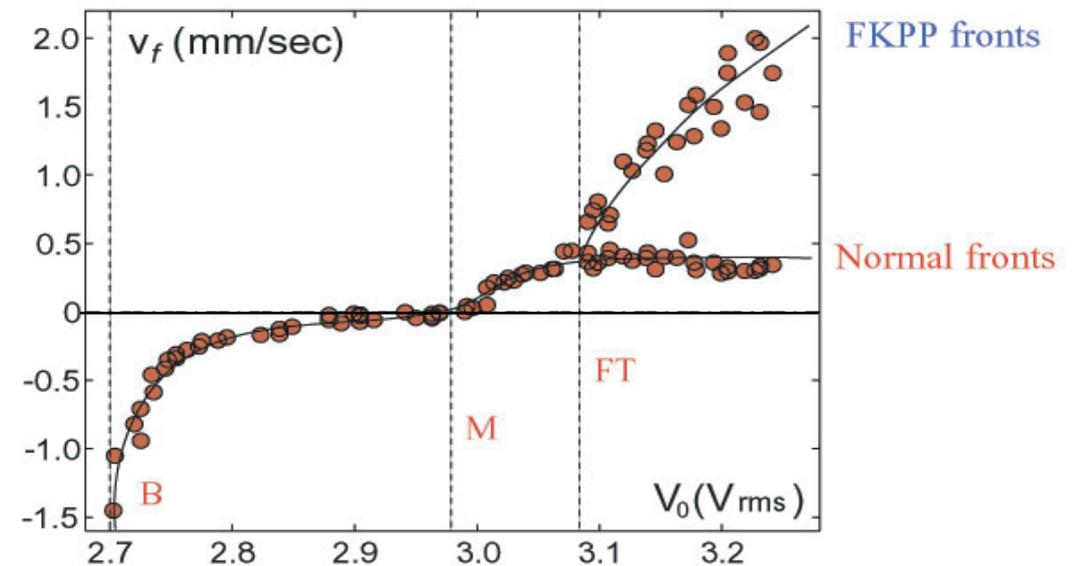
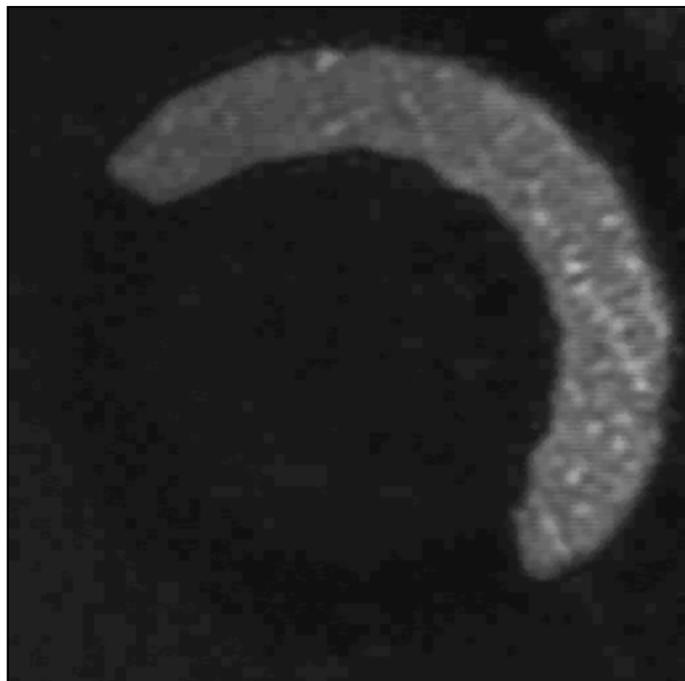
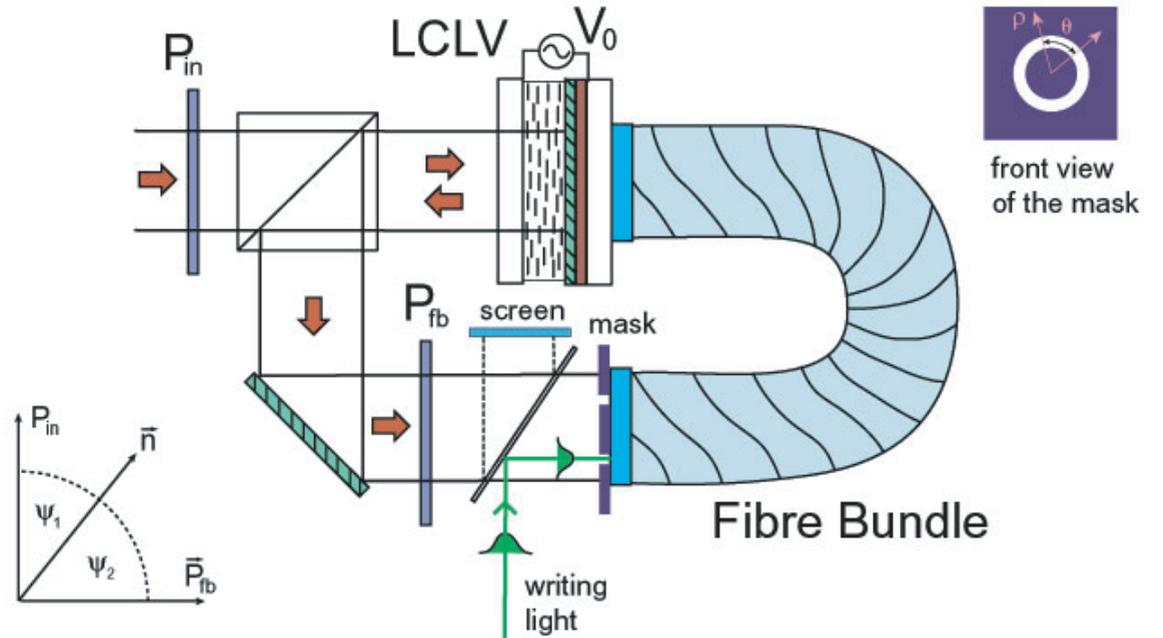
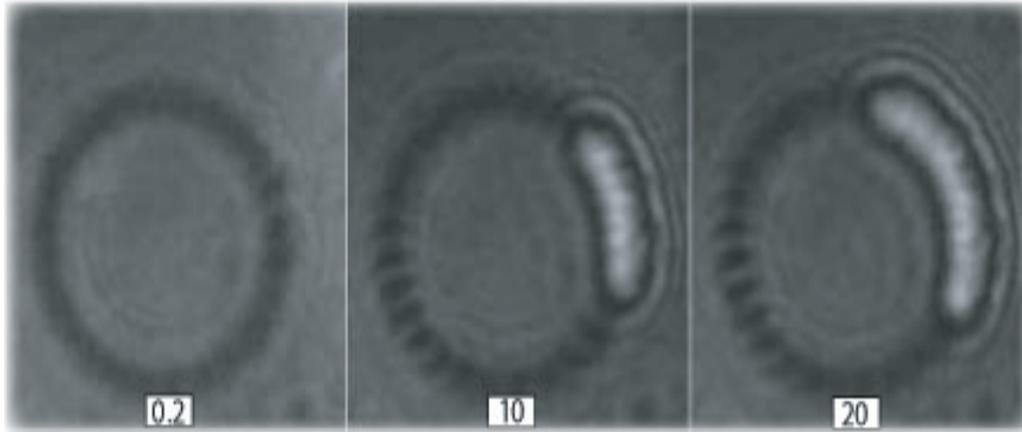


Pattern and uniform state



Front and experiments

- Liquid crystal light valve with optical feedback (1-D experiment)

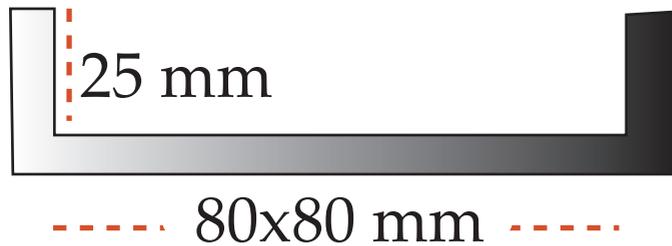


M.G. Clerc et al, Eur. Phys. J. D 28, 435 (2004).

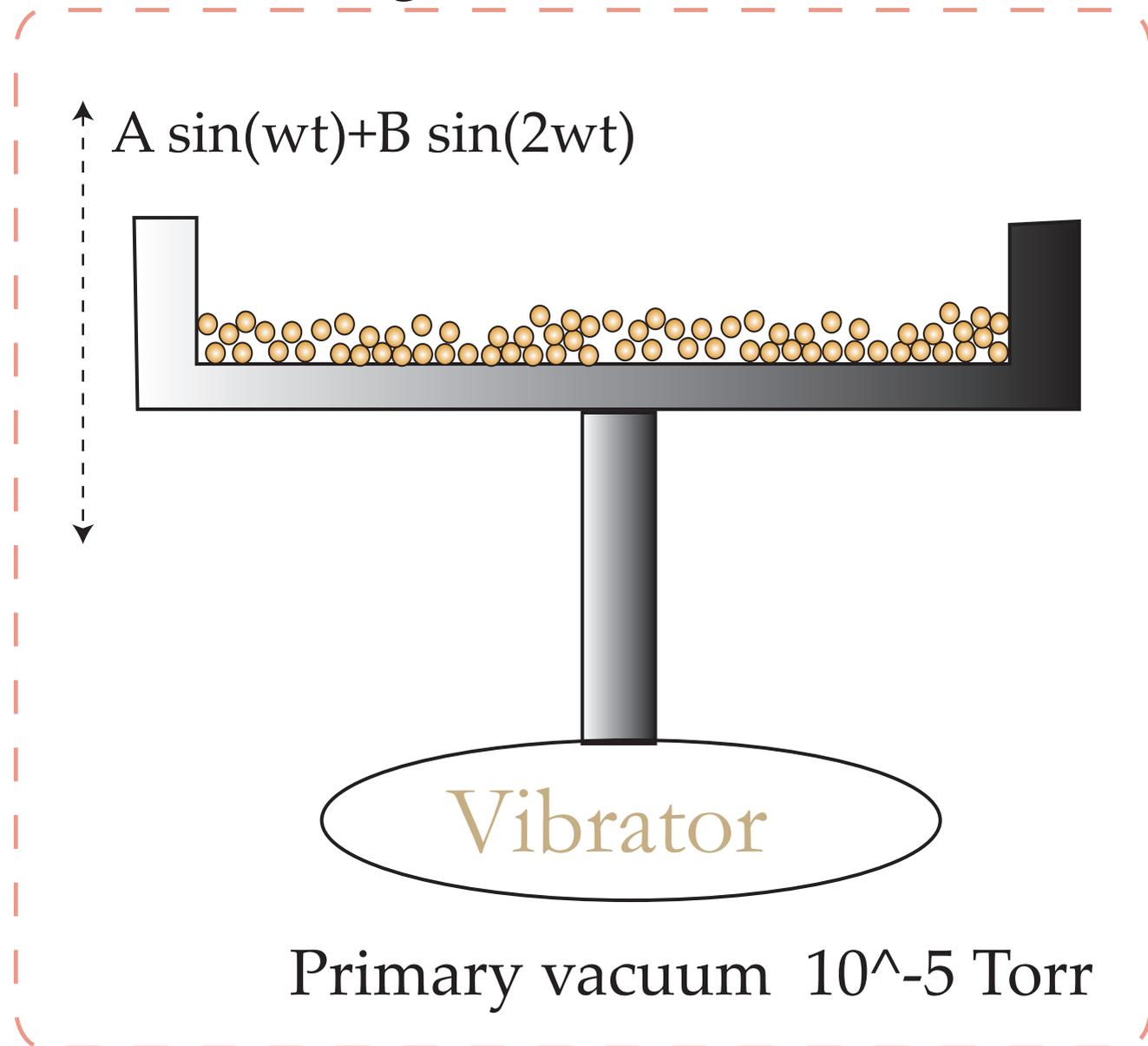
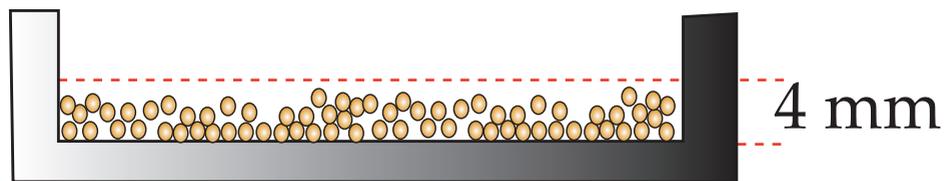
Pattern and uniform states in fluidized granular system

Experimental setup: container with broze grain

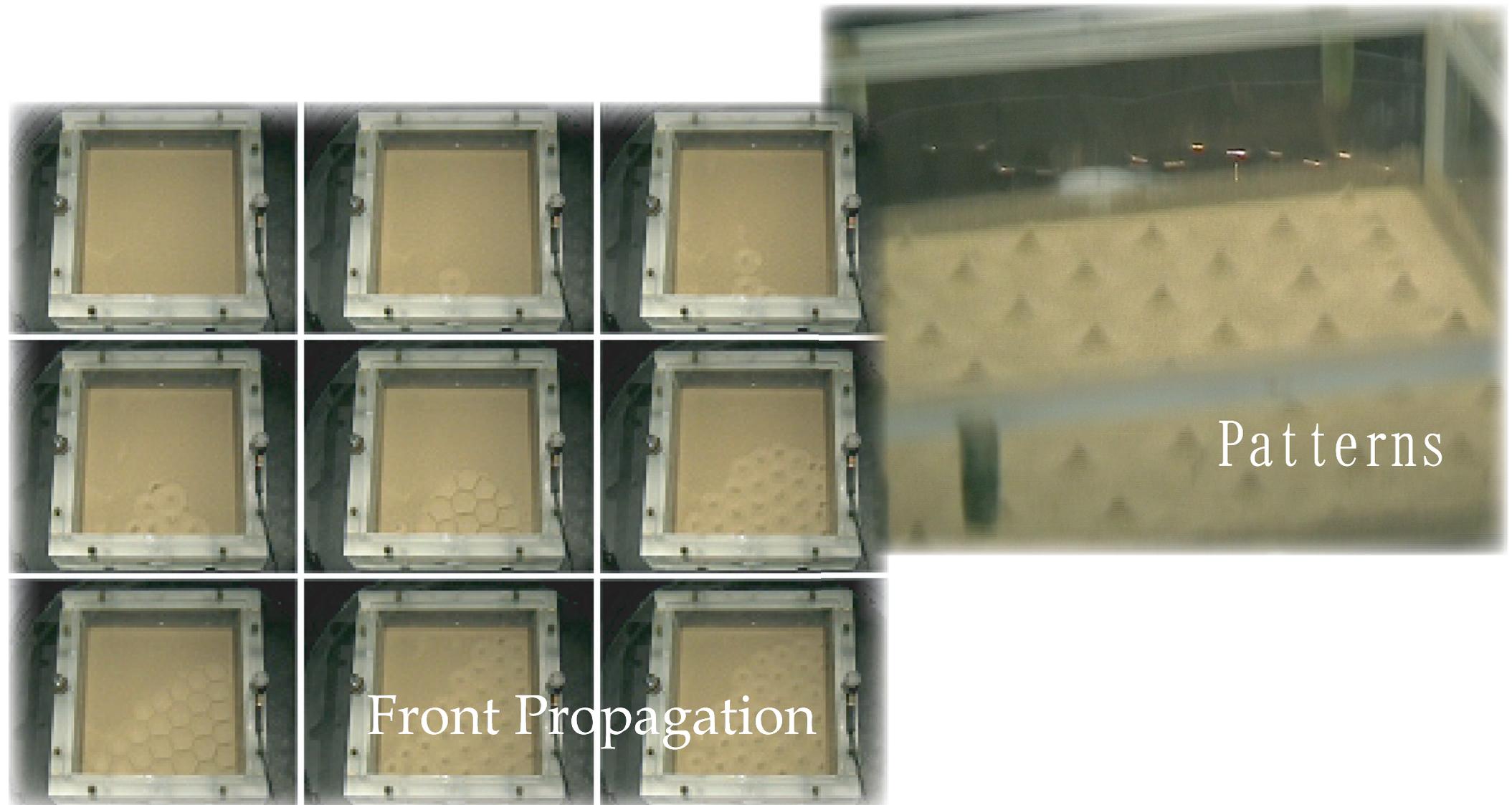
● 100 microns (grains)



Initial conditions



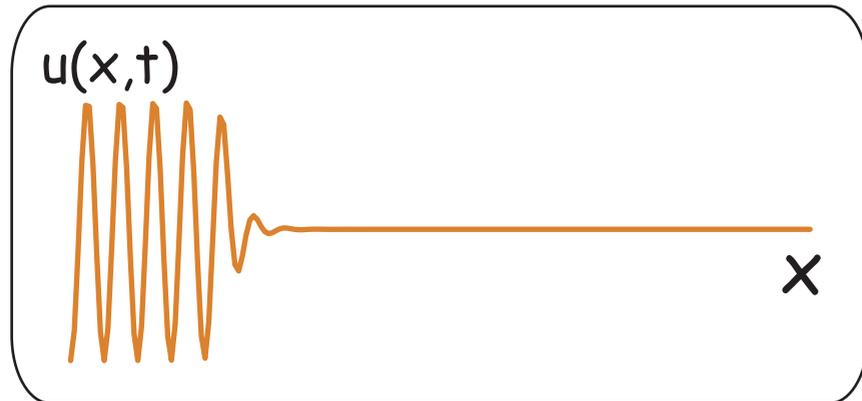
Front solutions



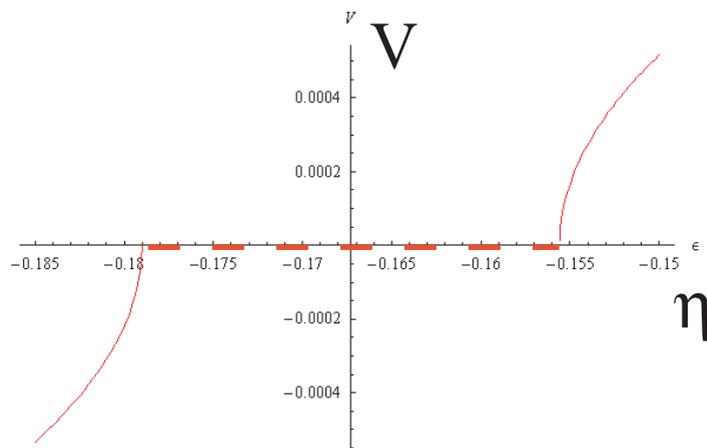
Courtesy of S. Residori, C. Falcon, U. Bortolozzo and M.G. Clerc (INLN)

Properties of front in one extended systems (1D)

- Locking phenomenon and pinning range

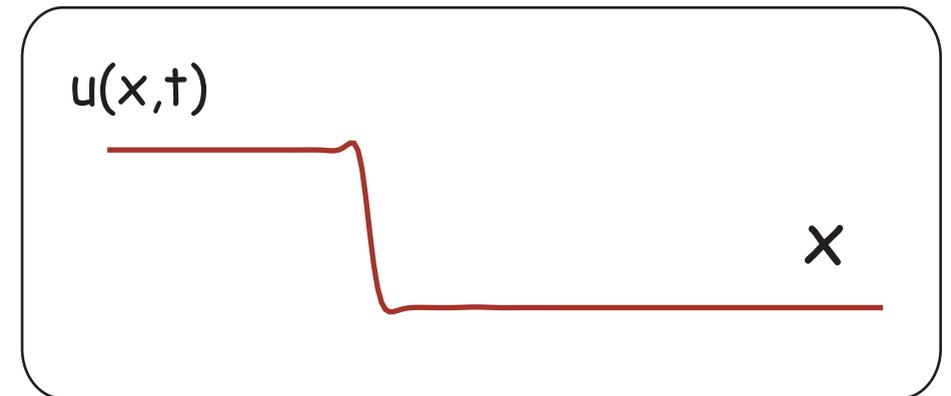


The front is stationary in a width range of parameters, pinning range

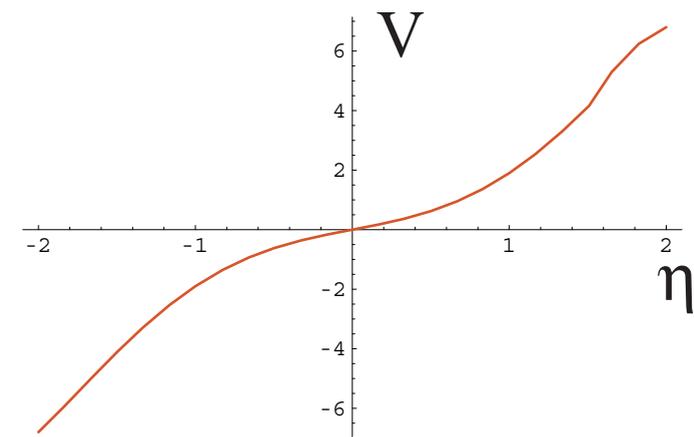


Y Pomeau, Physica D 23, 3 (1986)

- Normal form



The front is stationary in one point, **Maxwell point**.

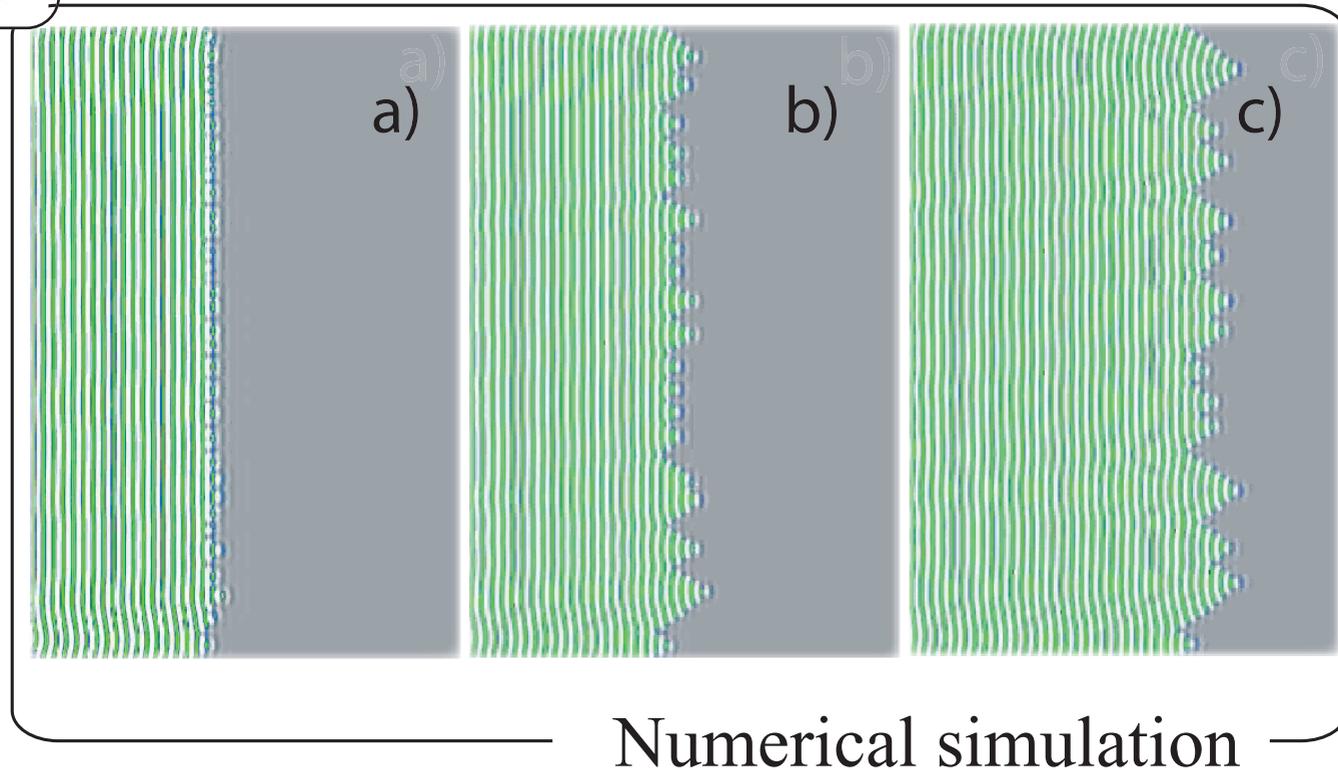
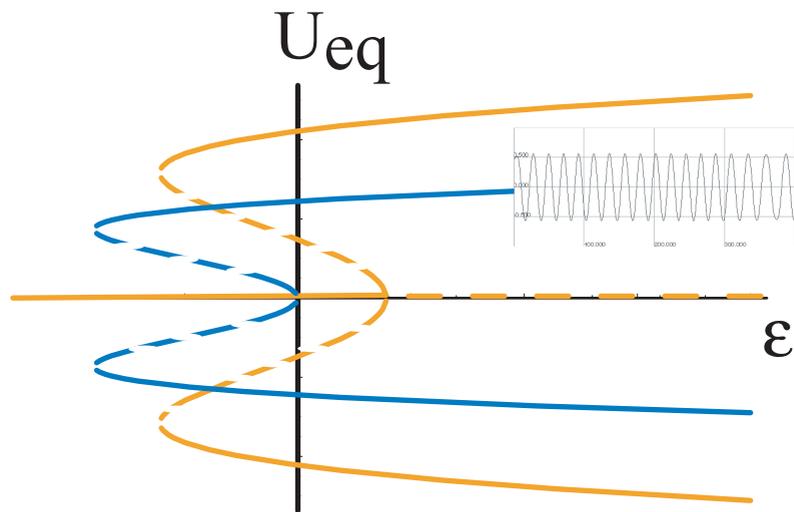


Which are the features of front in 2D?

- For the sake of simplicity, we consider front connecting a **stripe pattern** with a uniform state. Numerical simulations of the *isotropic Swift-Hohenberg equation* show a complex dynamics of front solution.

$$\partial_t u = \varepsilon u + \nu u^3 - u^5 - (\nabla^2 + q^2)^2 u,$$

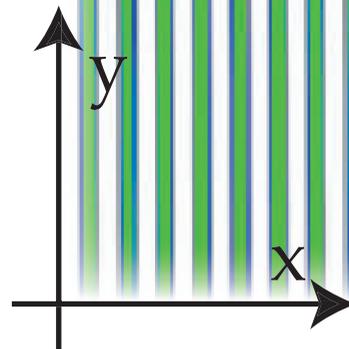
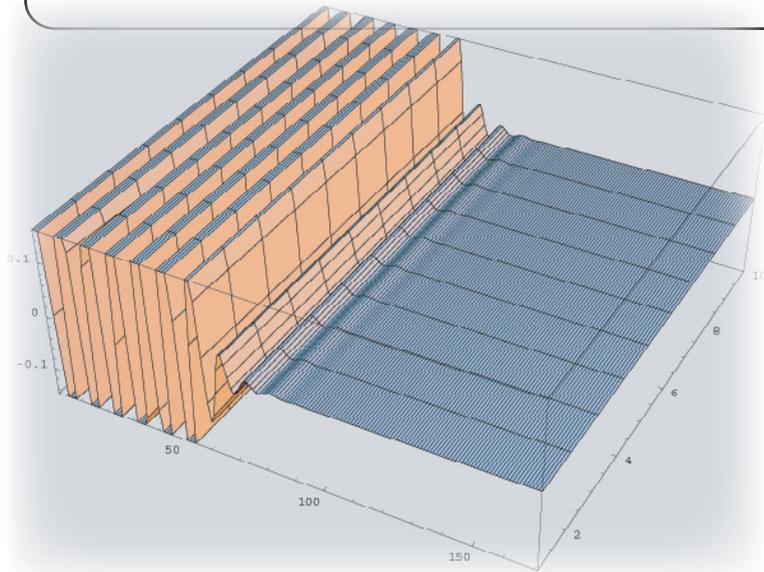
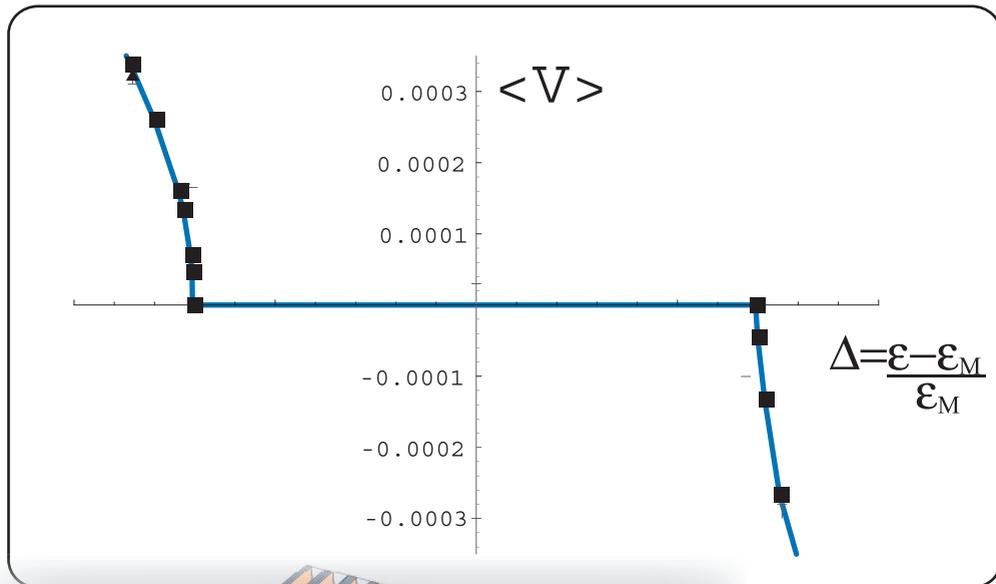
Bifurcation Diagram



Numerical simulation

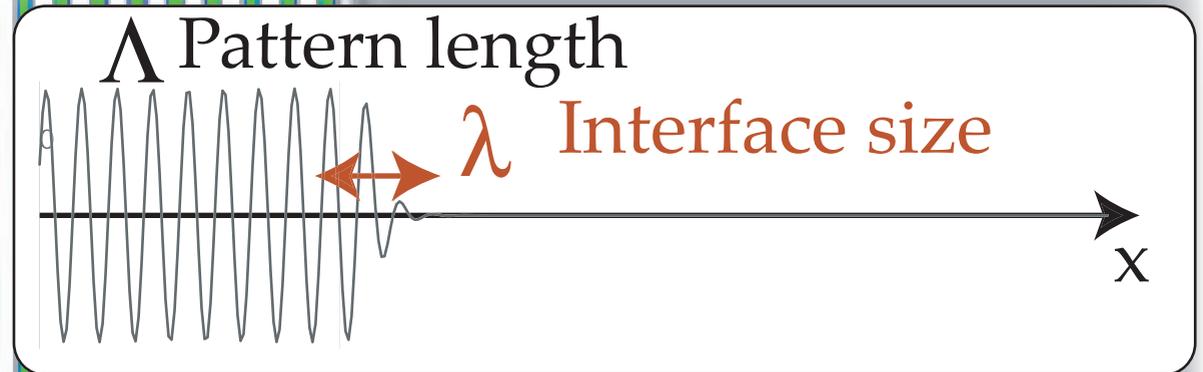
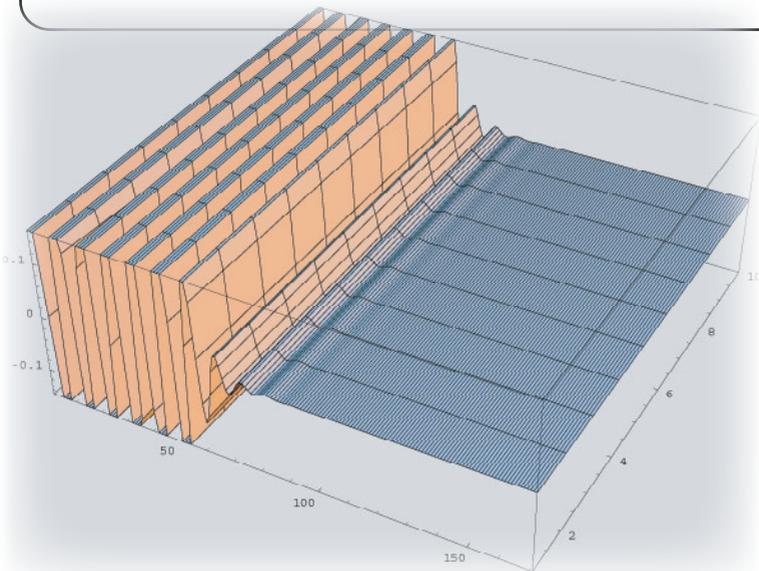
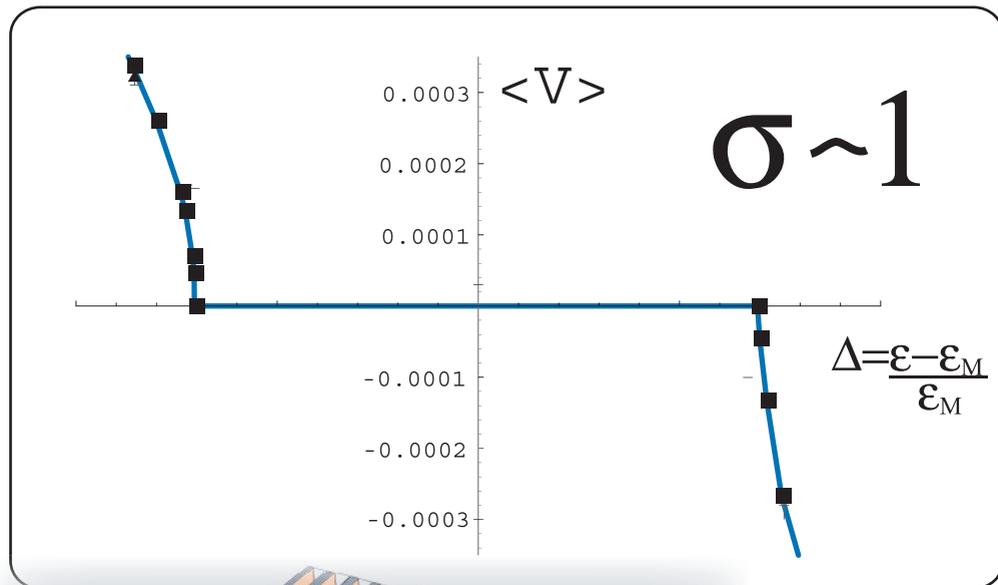
Locking phenomena

- The front speed is characterized by

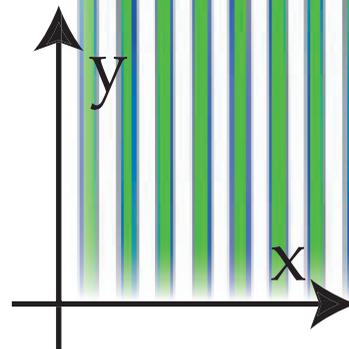


Locking phenomena

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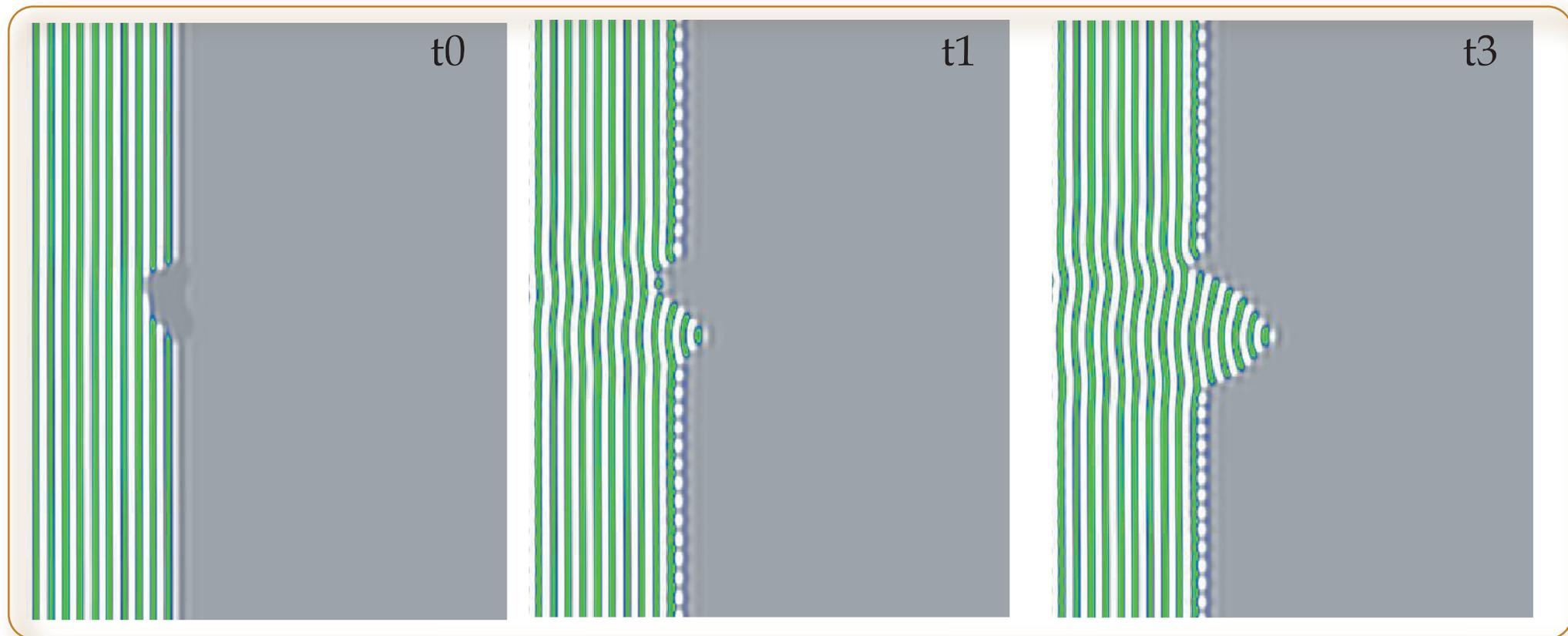


$$\sigma := \lambda / \Lambda$$



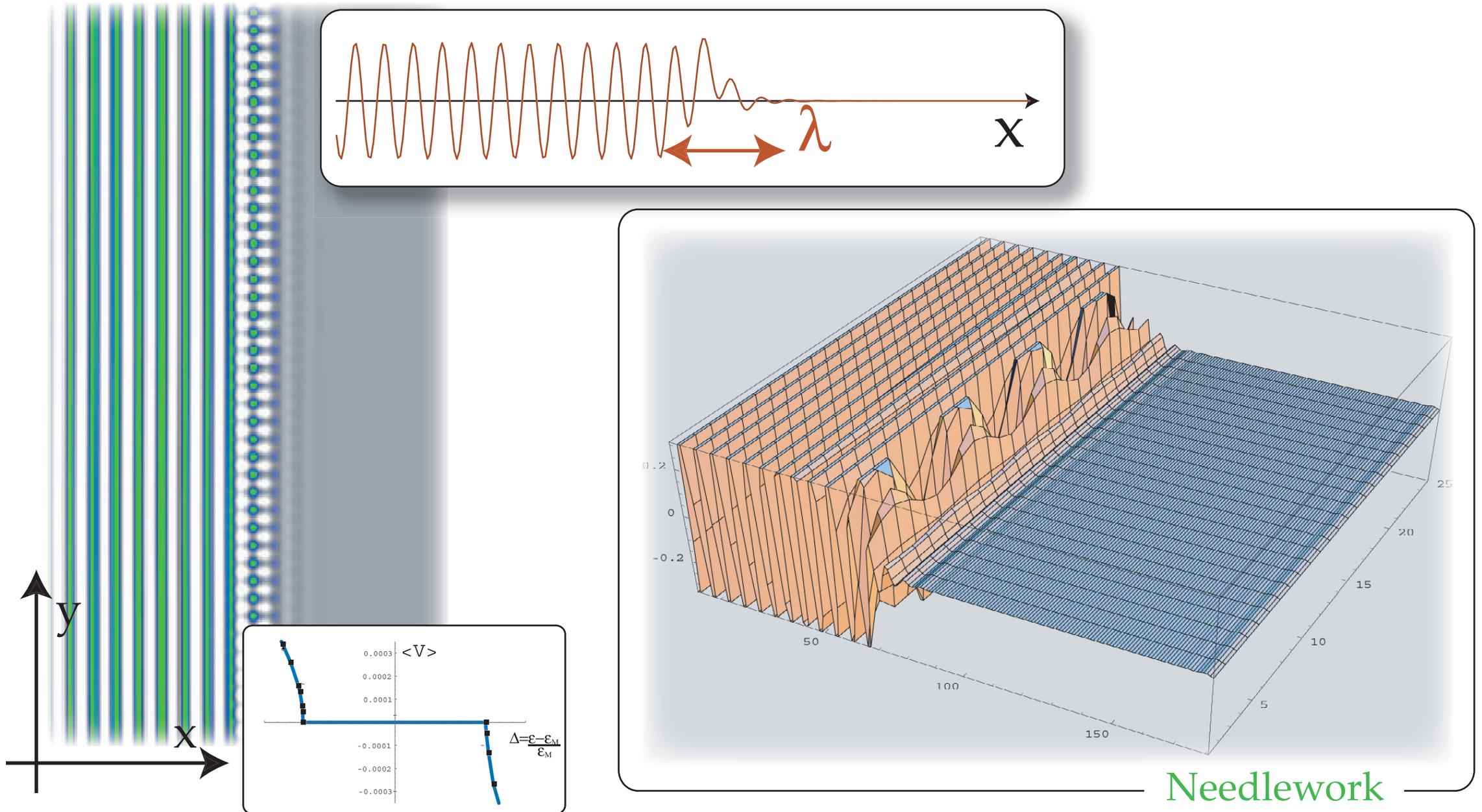
Zig-Zag dynamics

- For a finite perturbation, the interface exhibits zigzag dynamics, which is characterized by coarsening evolution



Needlework (Bordado, Broderie,..)

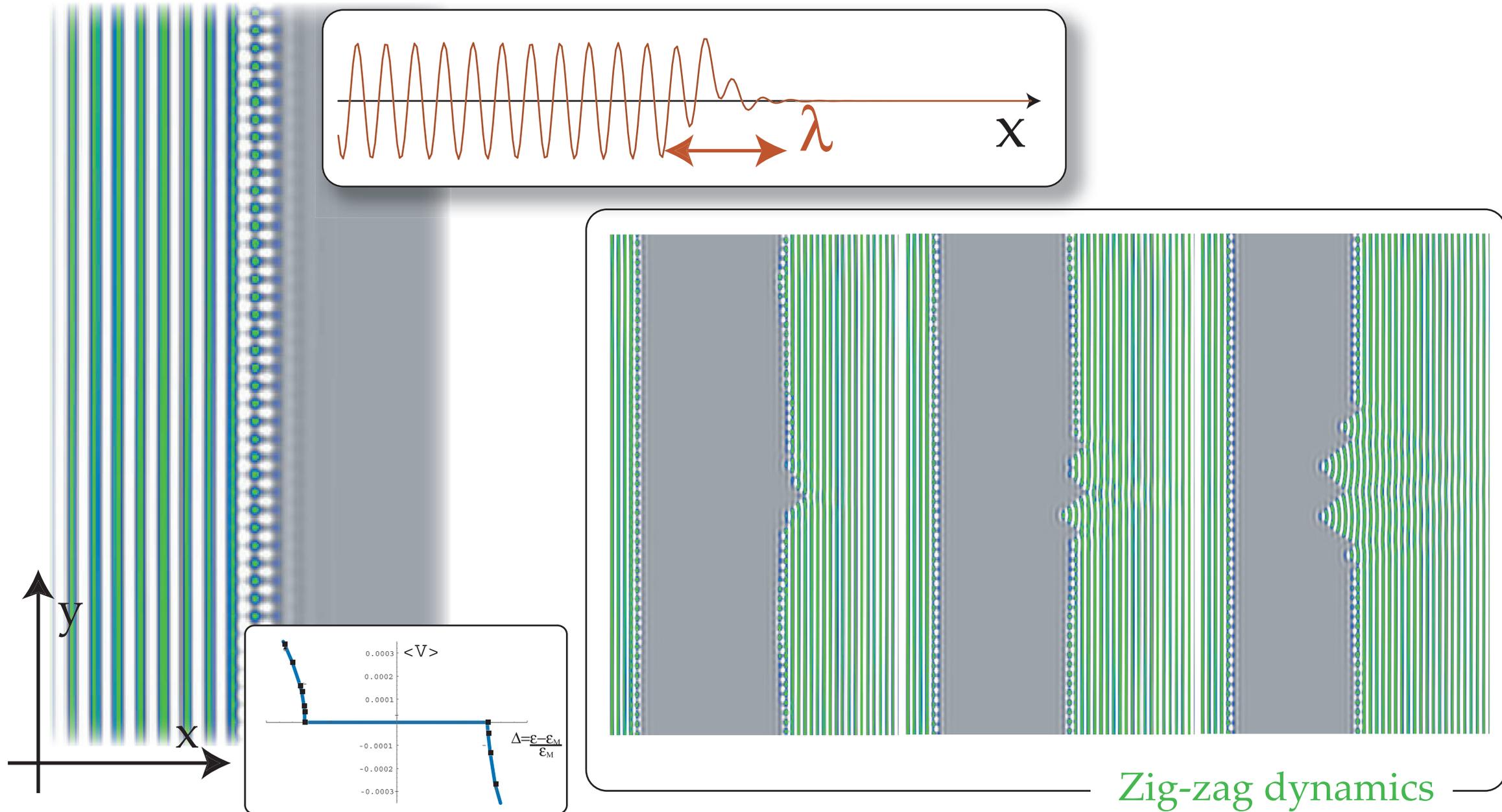
- For $\sigma > 1$, the system exhibits a spatial interface instability



Needlework

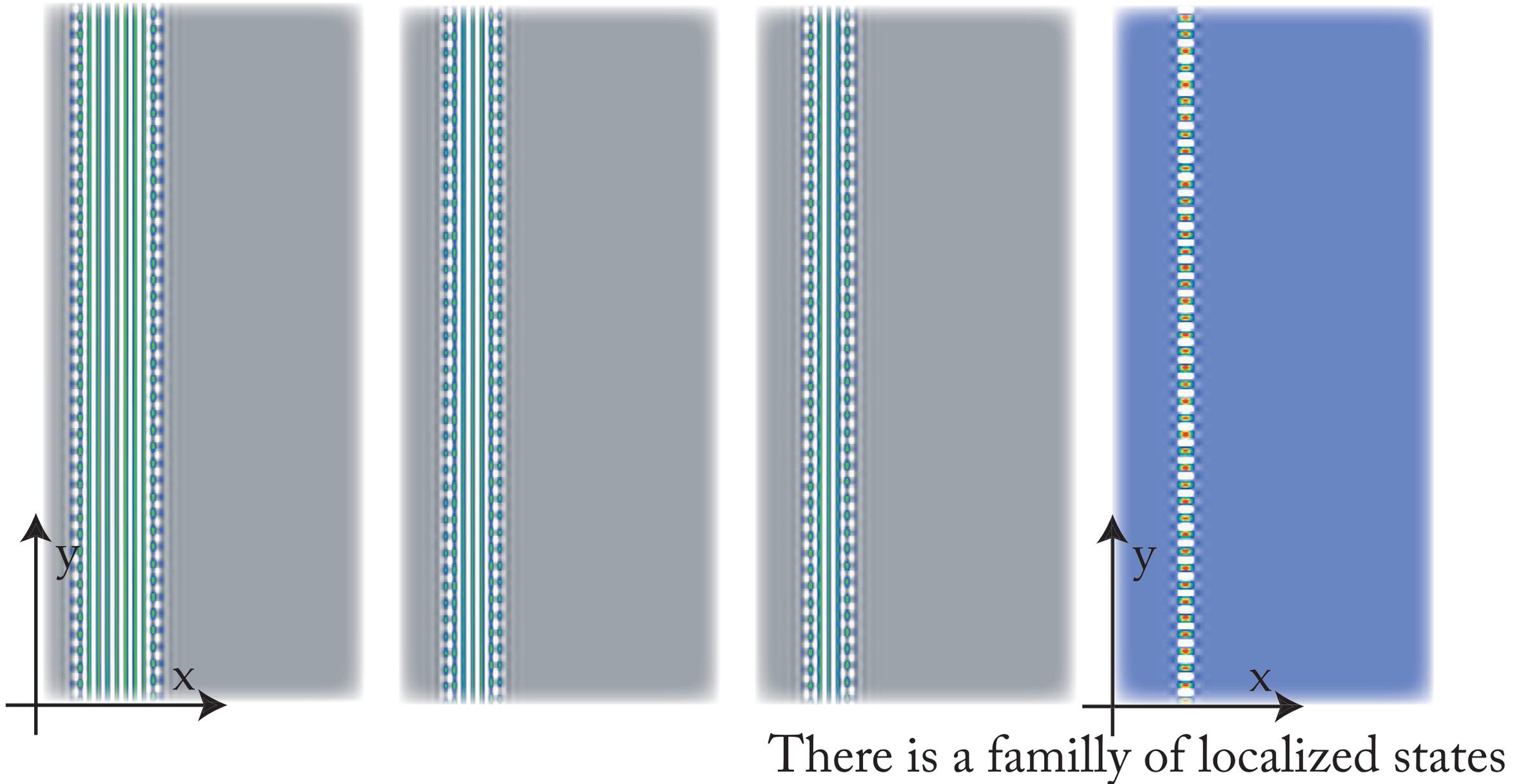
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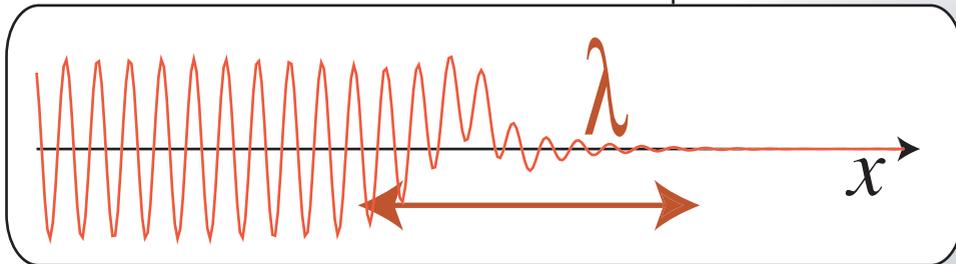
Localized states

- As consequences of locking phenomenon:

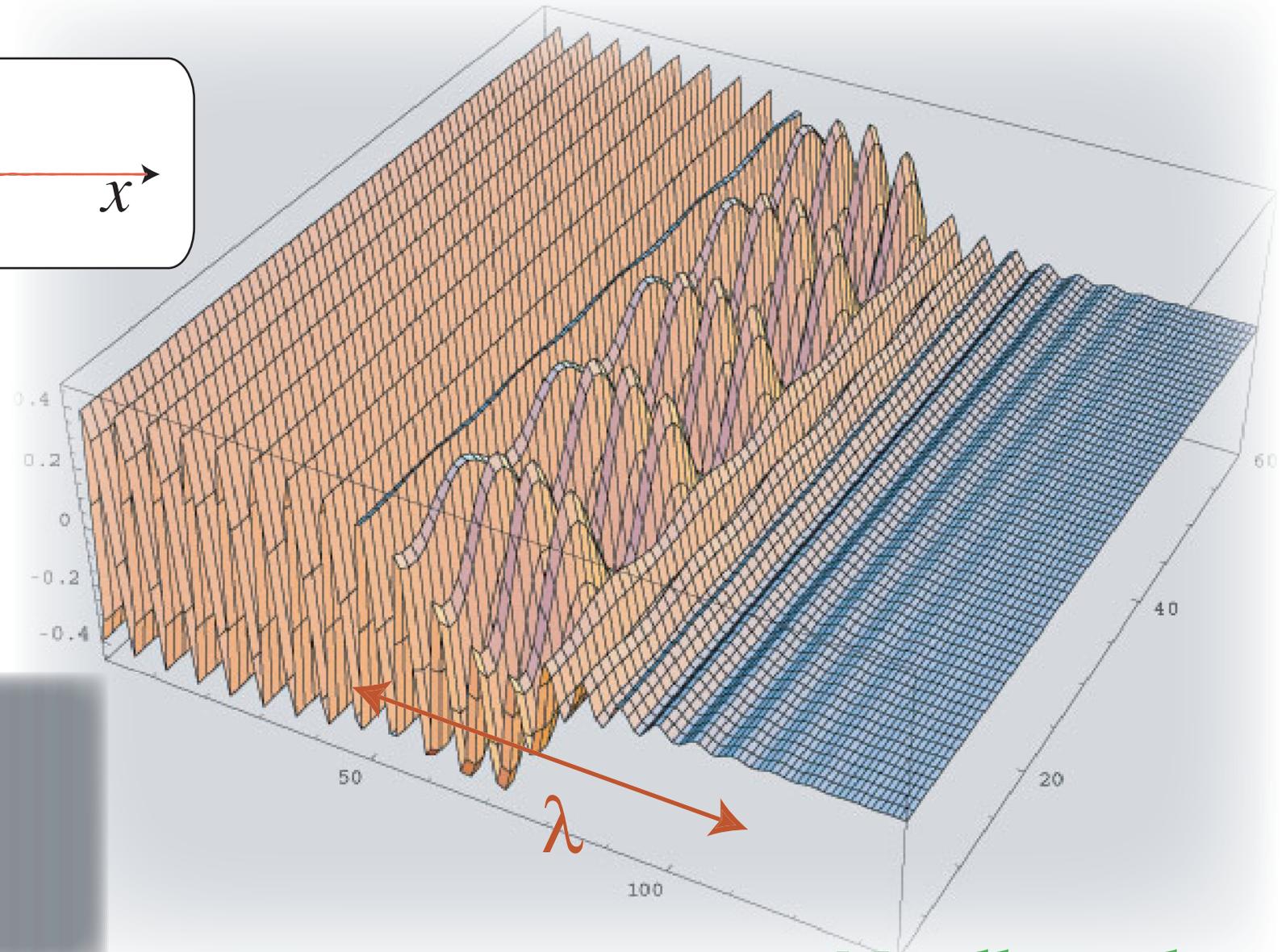
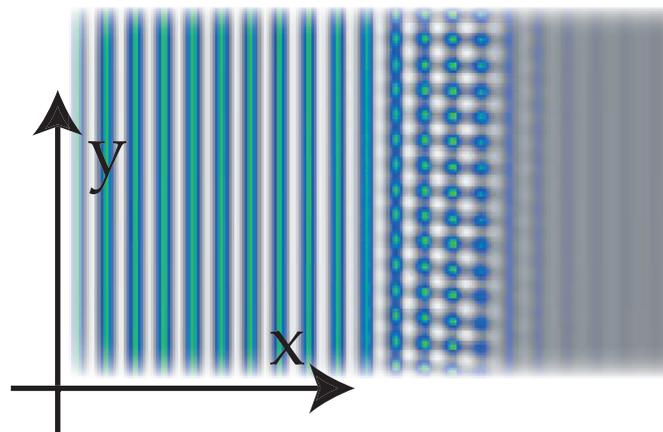


Needlework (Bordado, Broderie,..)

- Increasing $\sigma \gg 1$ (large σ), the needlework persists



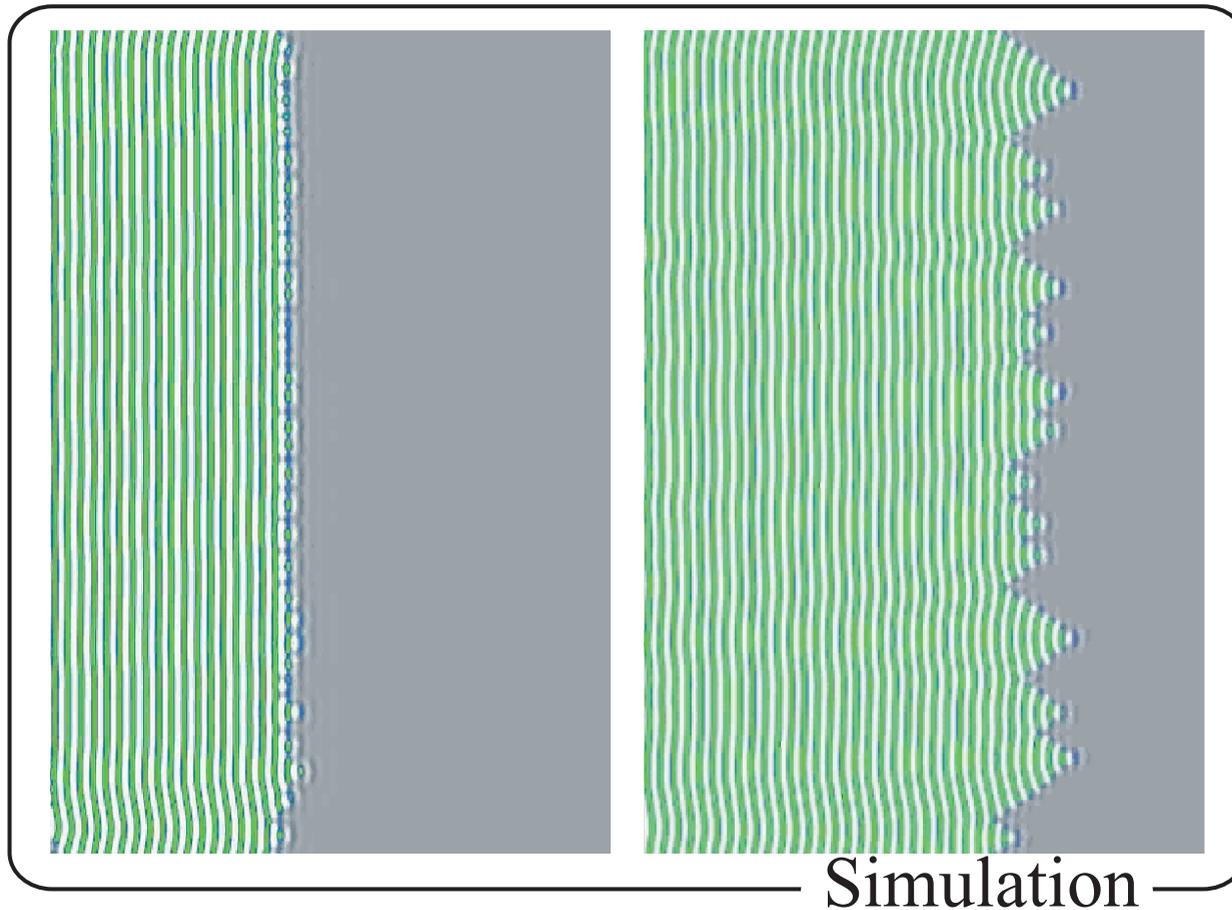
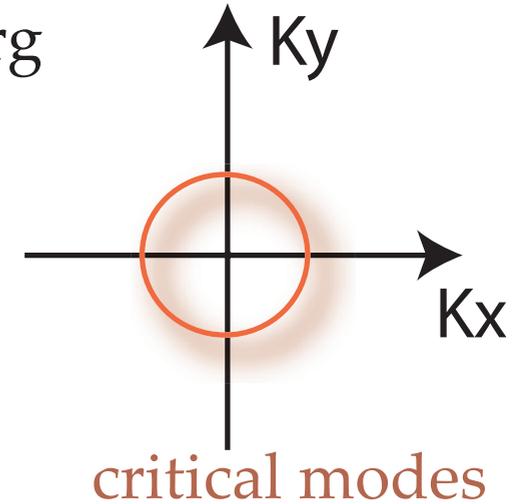
- The system has *pinning range* and it exhibits *zigzag dynamics*



Needlework

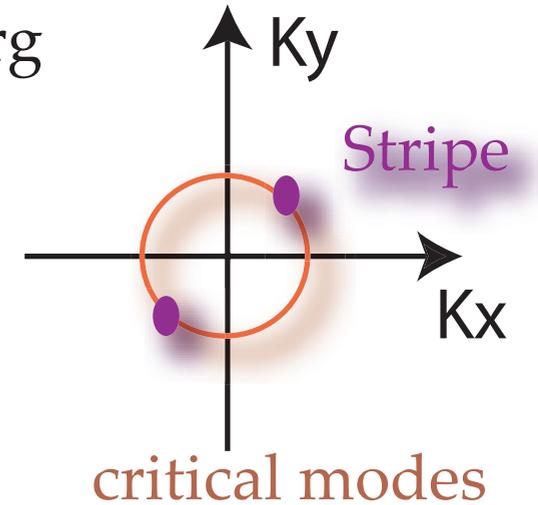
Mechanism of zigzag dynamics

- The zigzag dynamics exhibited by model Swift-Hohenberg model is triggered by the isotropy features of stripe orientation, that is, the orientation of the stripe is arbitrary and it is determined by the initial condition, whose influence is enhanced in the wall.

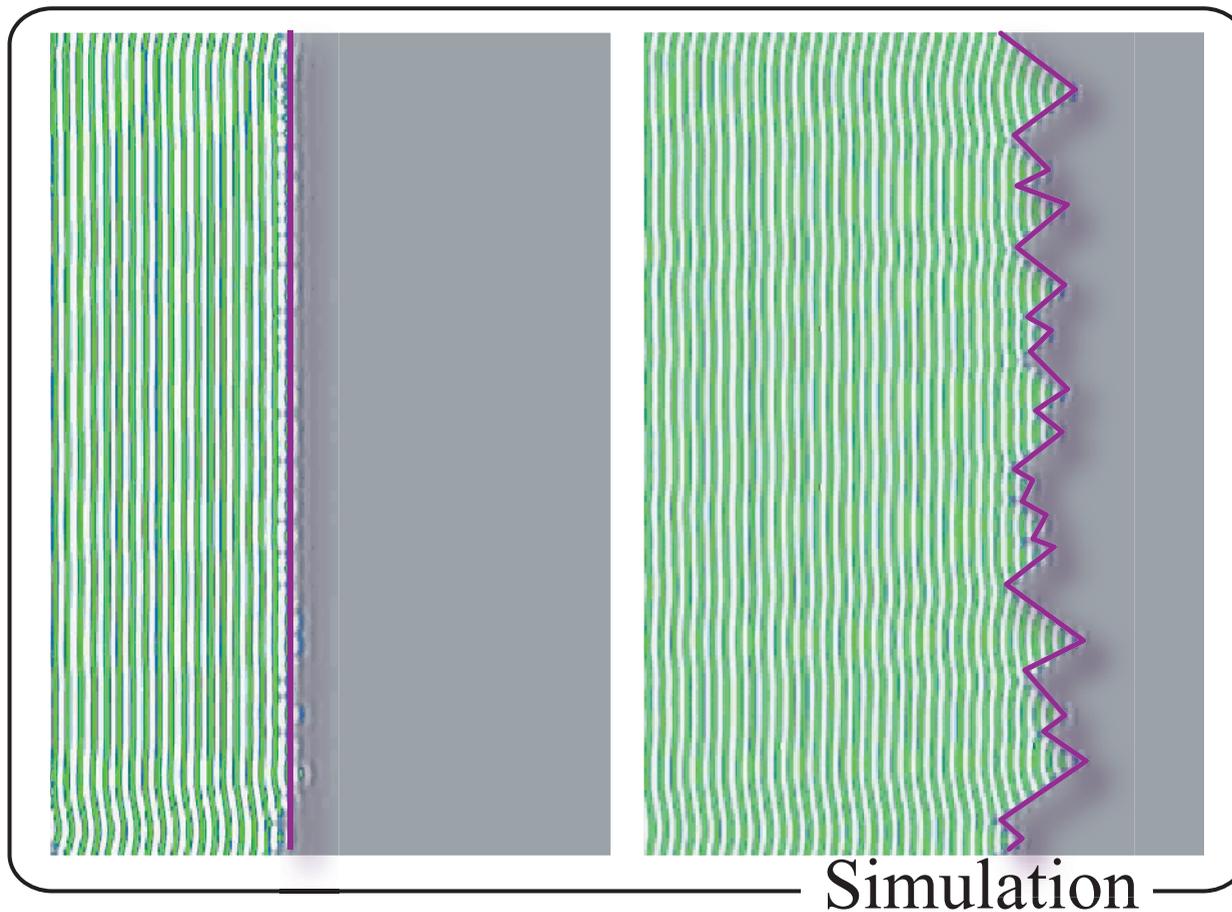


Mechanism of zigzag dynamics

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Zigzag dynamics and coarsening



Universal description

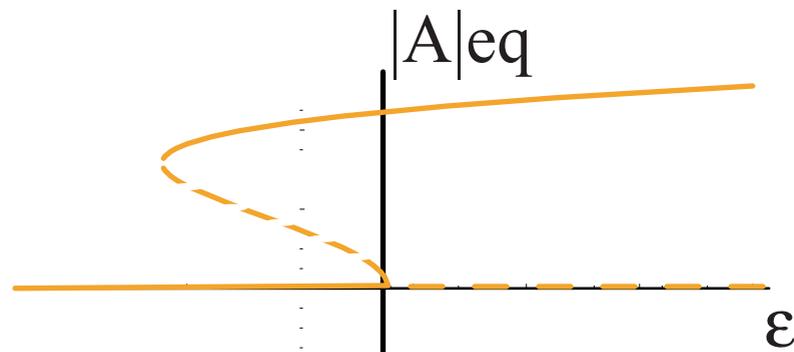
- For small ϵ and ν , we can introduce the ansatz

$$u = \sqrt{\frac{2\nu}{10}} \epsilon^{1/4} \left\{ A \left(X = \frac{x}{l_0}, Y = \frac{y}{l_0}, \tau = \frac{9\nu^2 |\epsilon|}{10} t \right) e^{iqx} + w_1(x, y, \tau) + c.c. \right\},$$

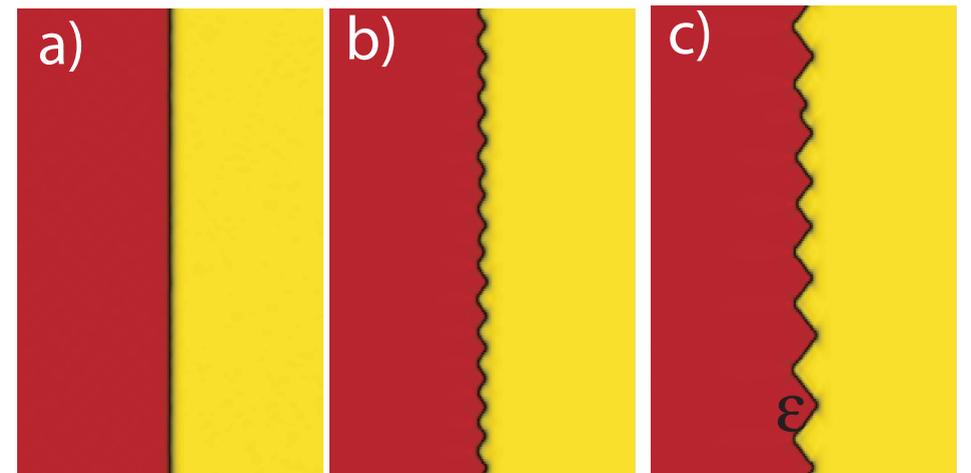
where the amplitud equation satisfies (Newell-Whitehead-Segel equation)

$$\partial_\tau A = \epsilon A + |A|^2 A - |A|^4 A + \left(\partial_x - i \frac{2}{q^2} \partial_{yy} \right)^2 A$$

Bifurcation Diagram



Numerical simulation



Universal description

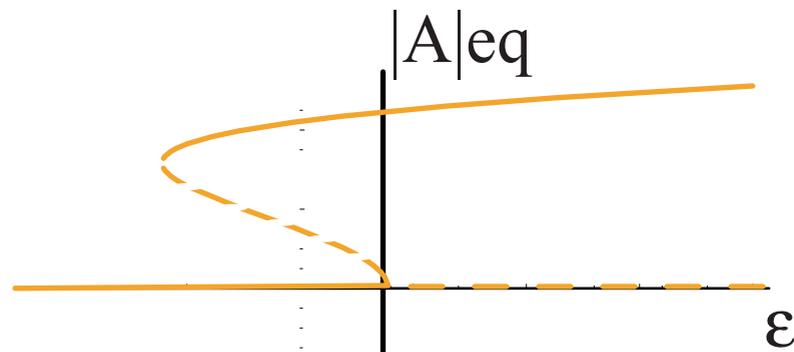
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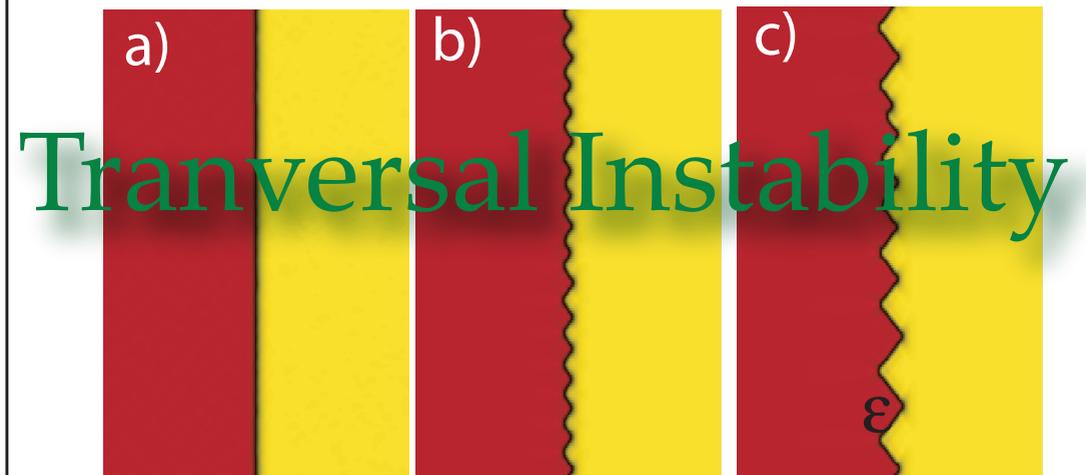
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Bifurcation Diagram



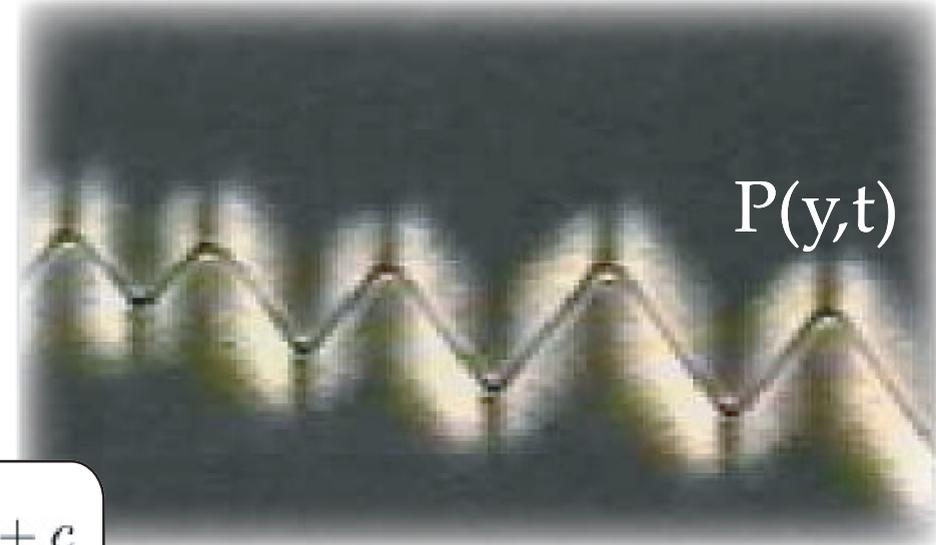
Numerical simulation



Interface equation

- A standard method to figure out the dynamics exhibited by the system is to derive an equation for the interface. The propagative zigzag interface is universally described by convective Cahn-Hilliard model

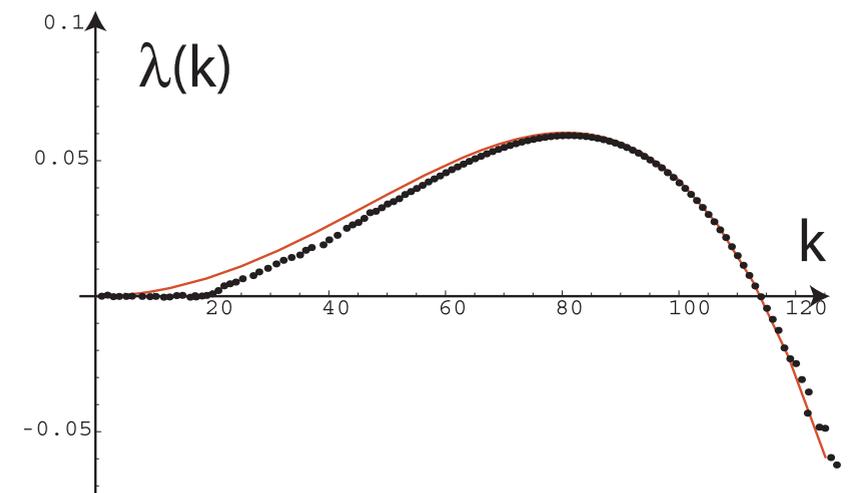
$$\partial_t P = \varepsilon P_{yy} + P_y^2 P_{yy} - P_{yyyy} + \alpha P_y^2 + b P_{yy}^2 + e P_y P_{yyy} + c$$



M. Clerc et al Eur. Phys. J. E 1, 179 (2000).

- Using this method, we have obtained a generalized Cahn-Hilliard equation, which shows that the flat interface is marginal (linearly) and nonlinear stable. Hence, this method does not give account of the depinning effect.

Spectrum



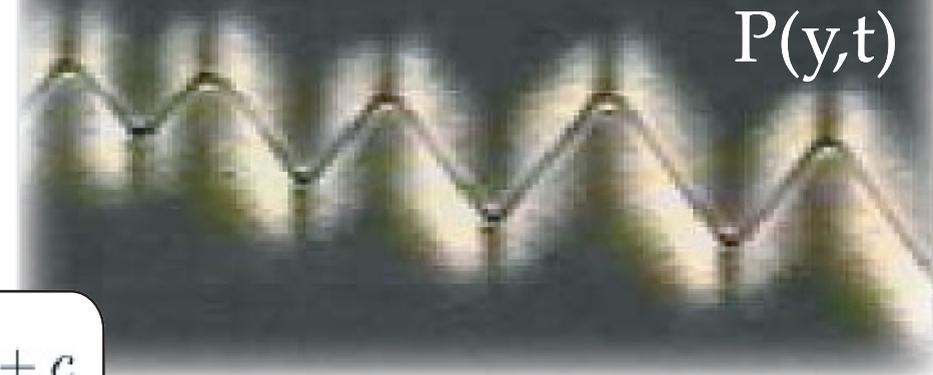
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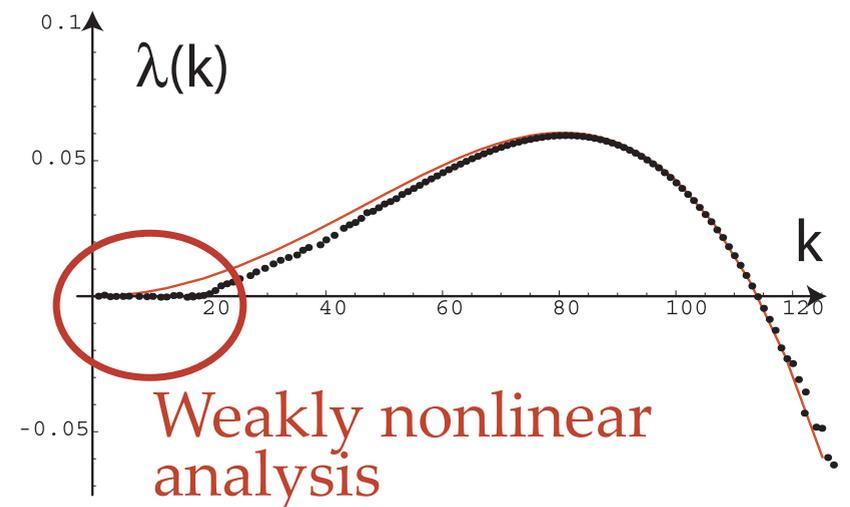
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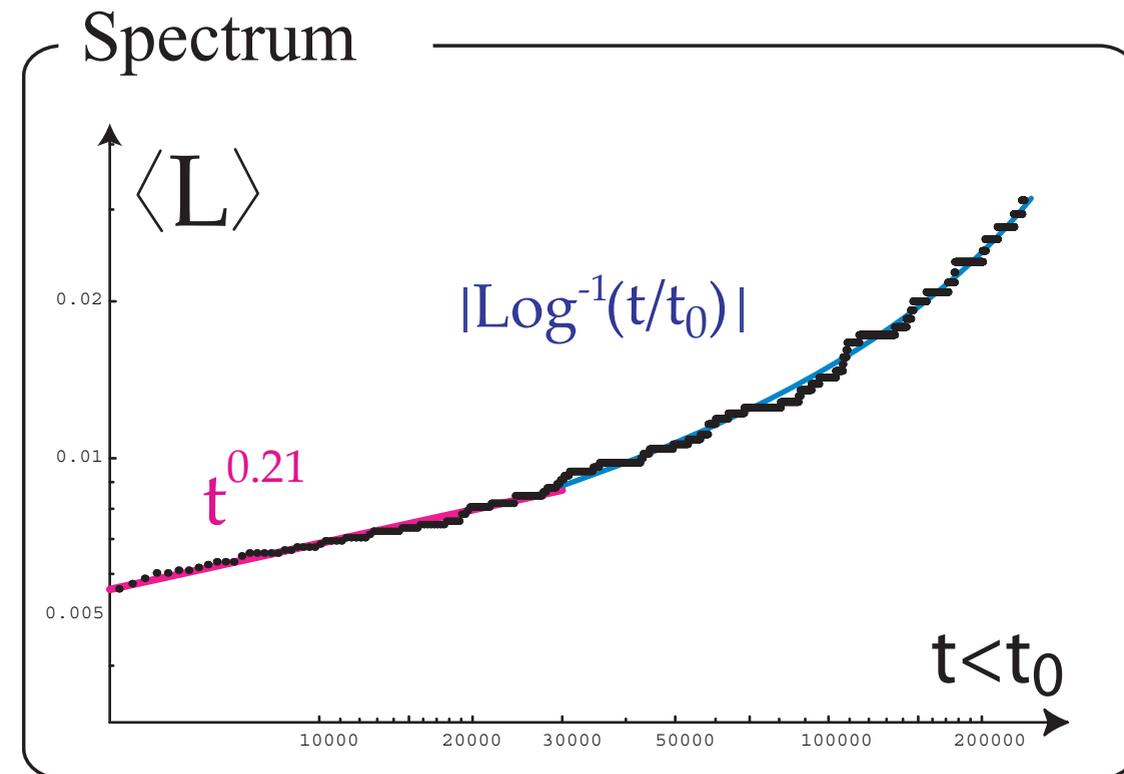
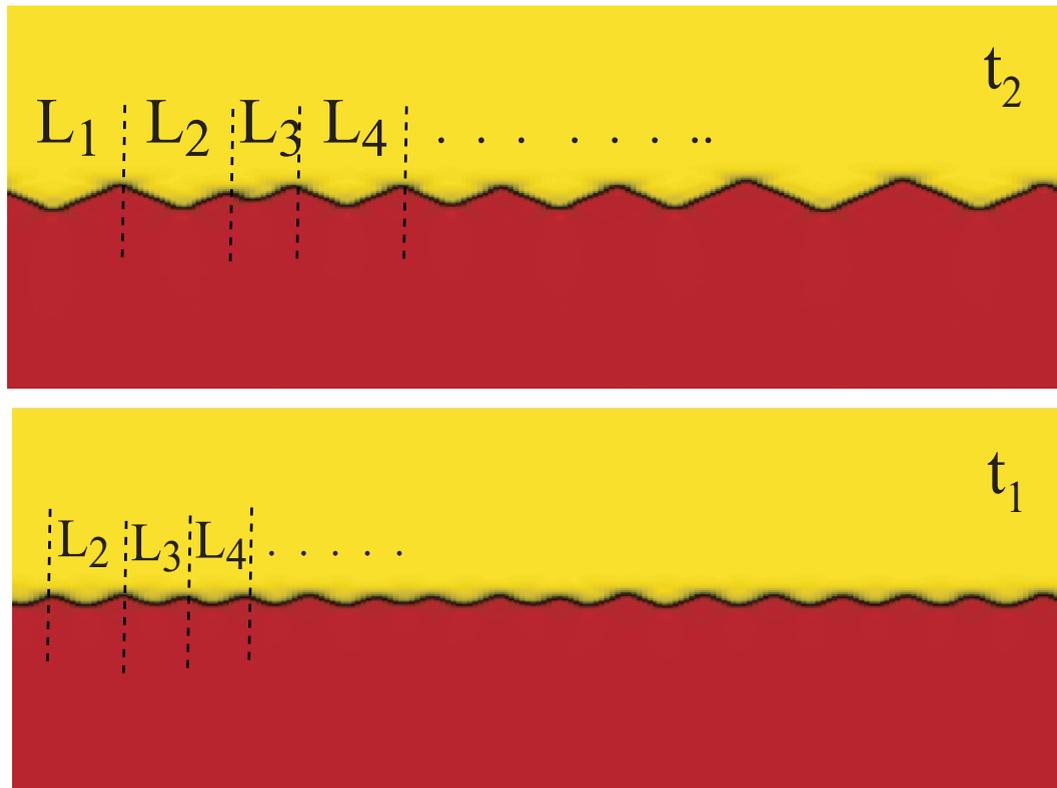


Spectrum



Coarsening

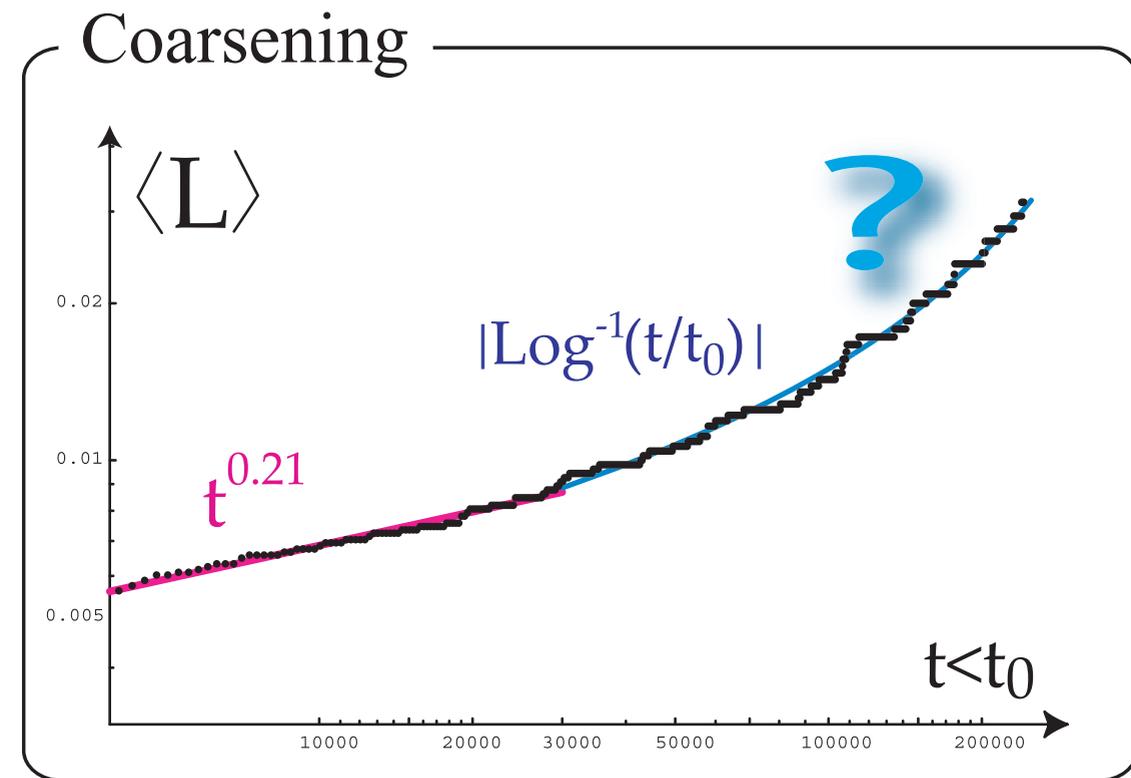
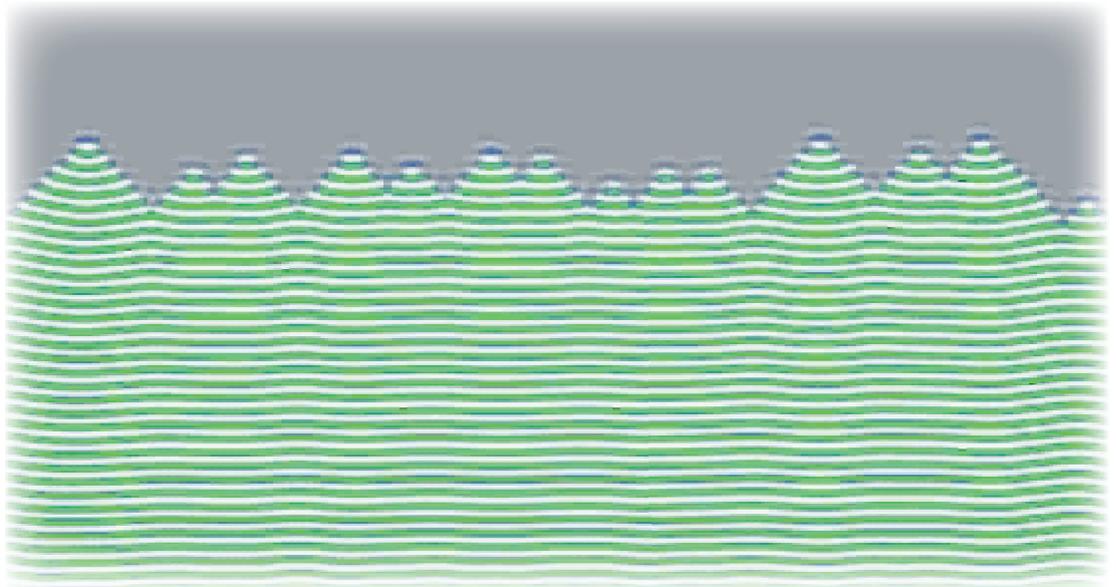
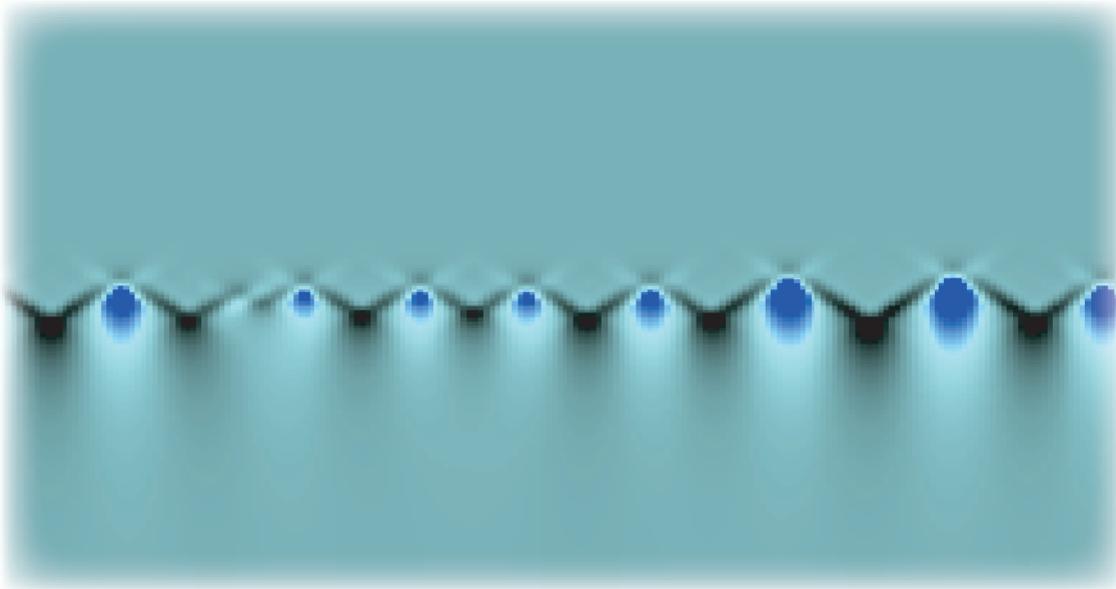
- Numerically, we compute the average length size $\langle L(t) \rangle$ between two successive extreme points of the interface of model



$$\langle L \rangle = \frac{1}{N} \sum_{i=1}^N L_i$$

Coarsening

- Numerically, we compute the average length size $\langle L(t) \rangle$ between two successive extreme points of the interface of model



Interface-phase-interface interaction

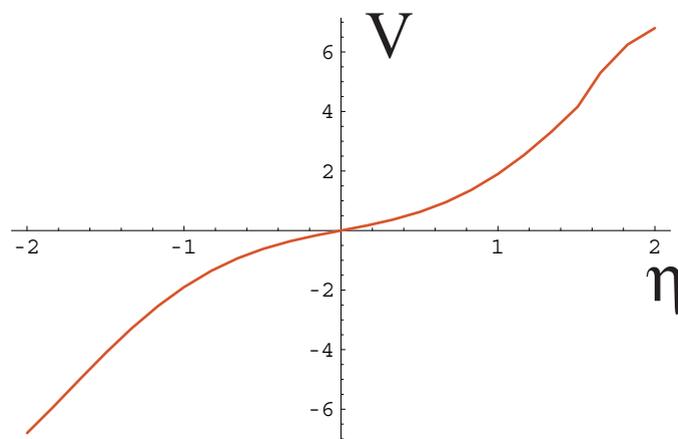
Amended amplitude equation

- To explain the appearance of pinning range, needlework and complex dynamics, we consider

$$\partial_t A = \epsilon A + |A|^2 A - |A|^4 A + \left(\partial_x - \frac{i}{q} \partial_{yy}\right)^2 A$$

Resonant terms

Front speed



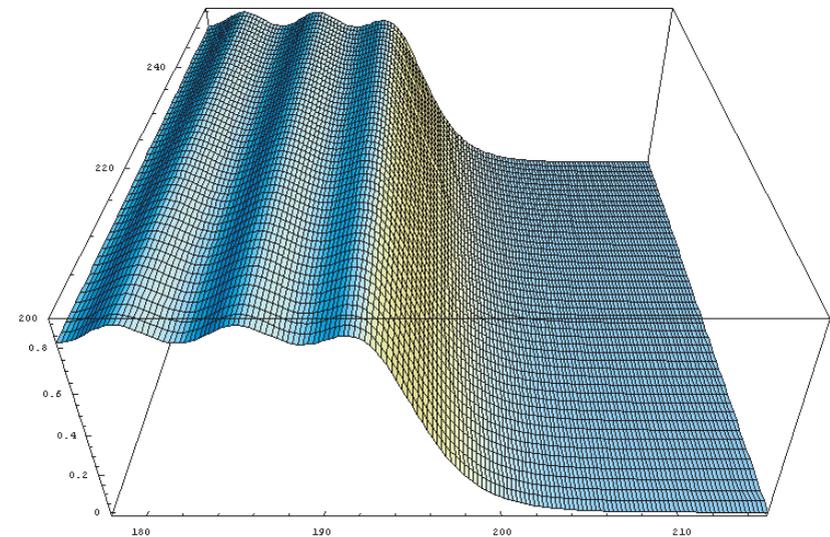
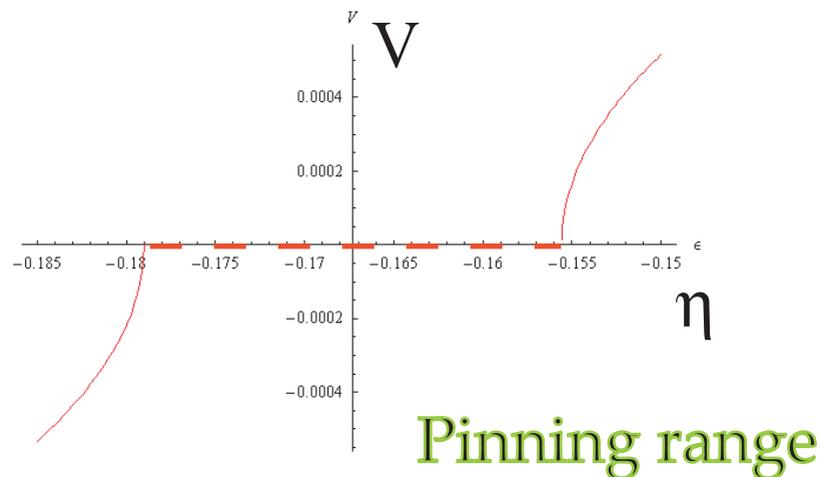
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Non Resonant terms

Front speed



large η

D. Bensimon, B.I. Shraiman, and V. Croquette, Phys. Rev.A 38, 5461 (1988).
M.G. Clerc, C. Falcon, and E. Tirapegui, Phys. Rev. Lett. 94, 148302 (2005).

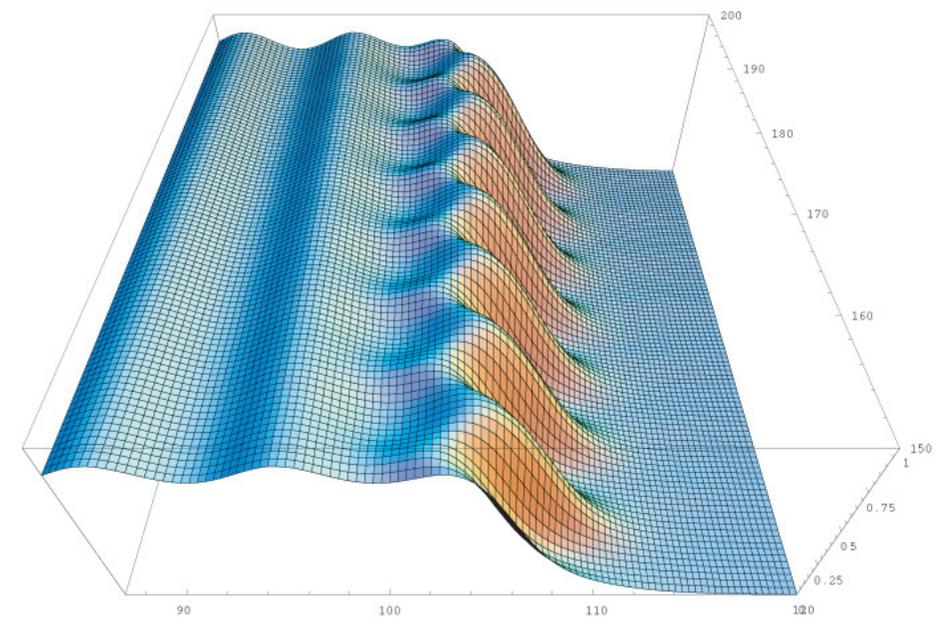
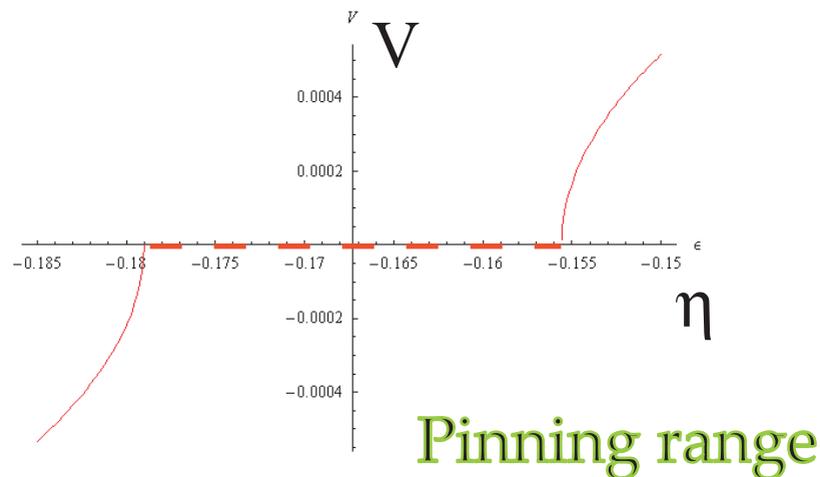
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Non Resonant terms

Front speed

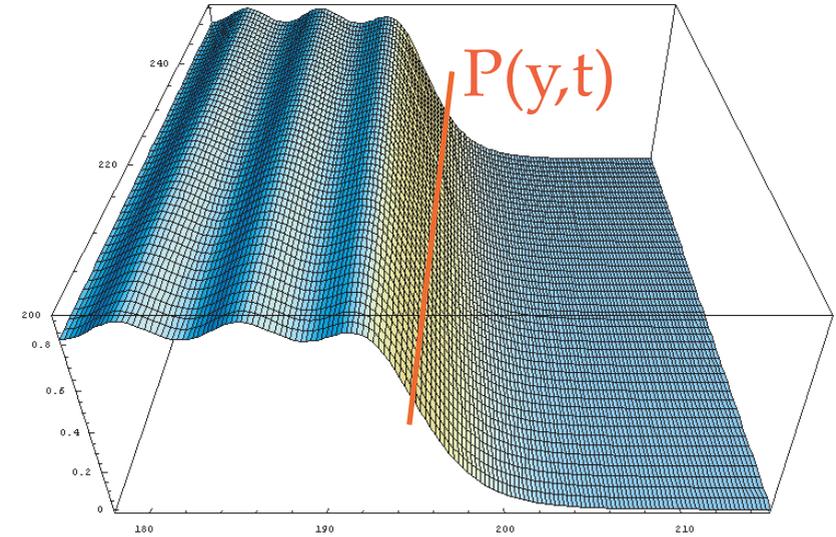


Smaller η

D. Bensimon, B.I. Shraiman, and V. Croquette, Phys. Rev.A 38, 5461 (1988).
 M.G. Clerc, C. Falcon, and E. Tirapegui, Phys. Rev. Lett. 94, 148302 (2005).

Simple prototype model

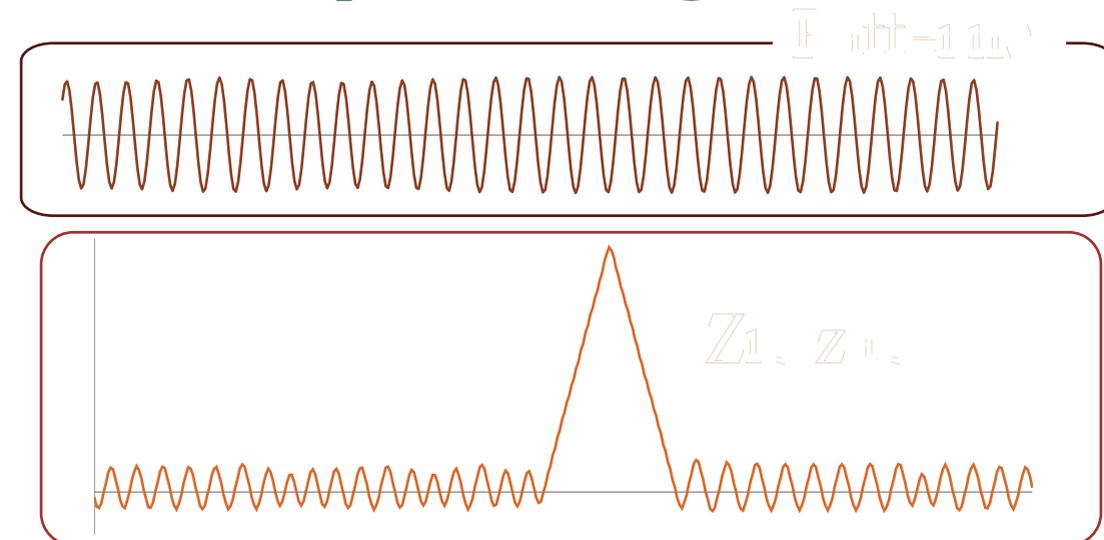
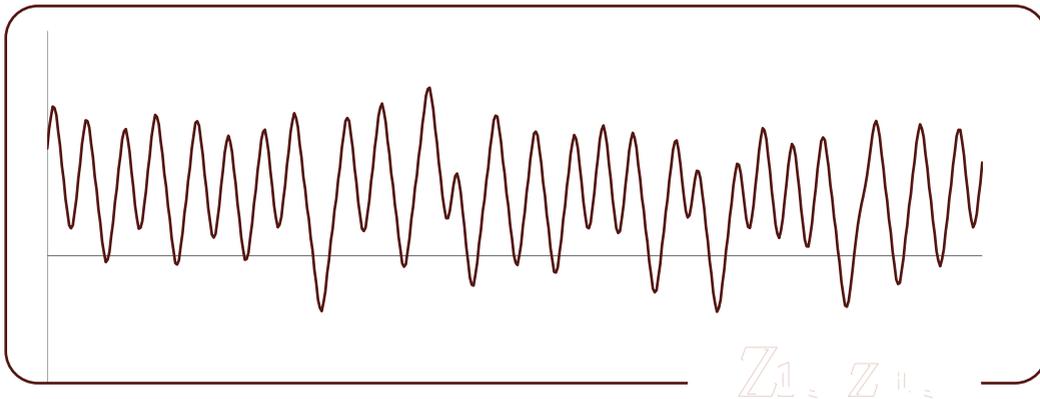
- $P(y,t)$ stand for the interface bewteen the stripe pattern and uniform state
- Modified Cahn-Hilliard equation



$$\partial_t P = \underbrace{\varepsilon P_{yy} + P_y^2 P_{yy} - P_{yyyyy}}_{\text{Zig-zag Instability}} + \underbrace{\alpha P_y^2 - \nu \sin(\kappa P)}_{\text{Spatial forcing}}$$

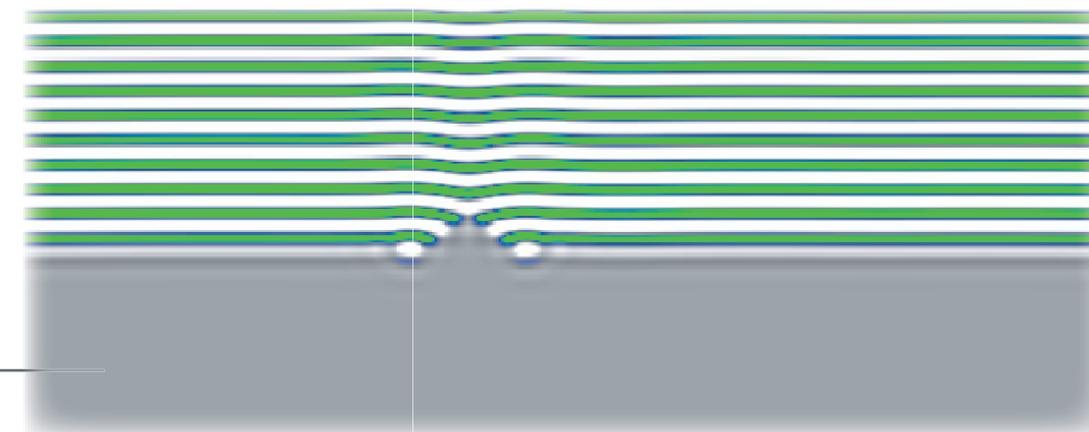
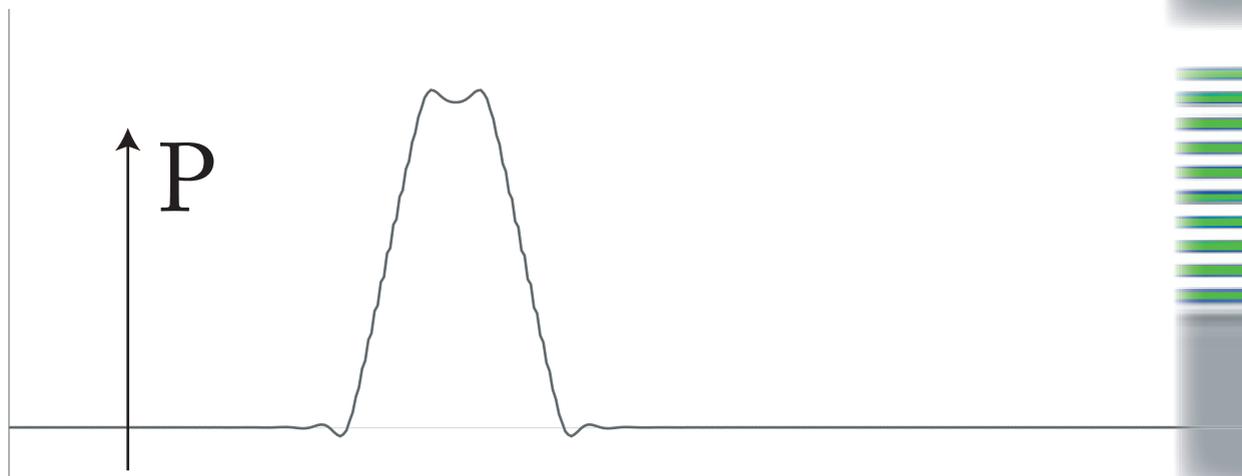
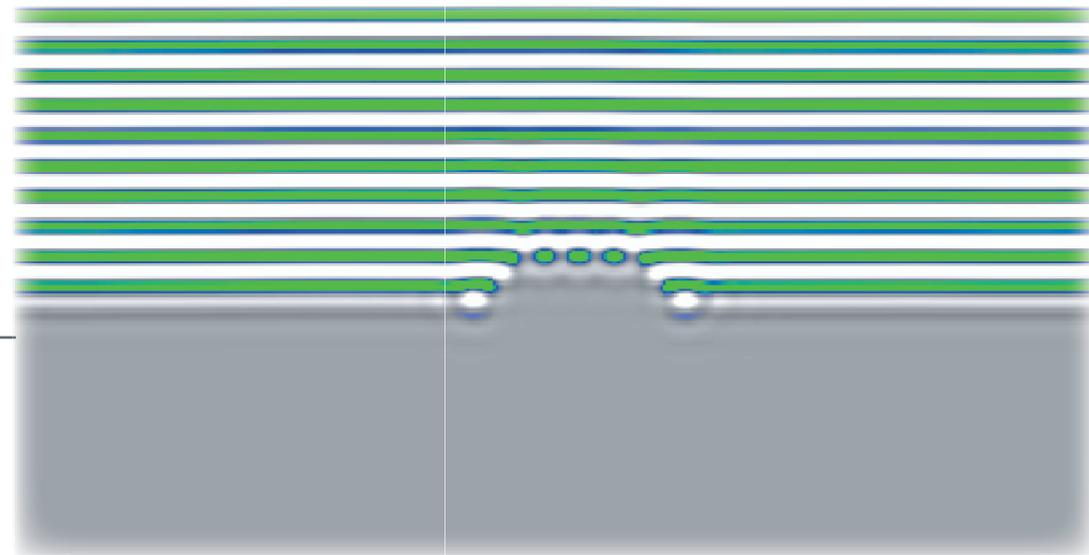
Zig-zag Instability

Spatial forcing



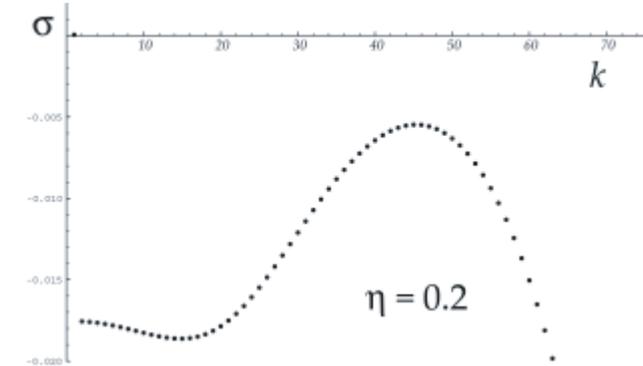
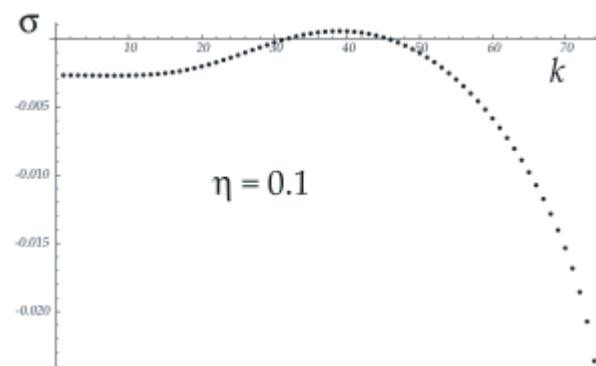
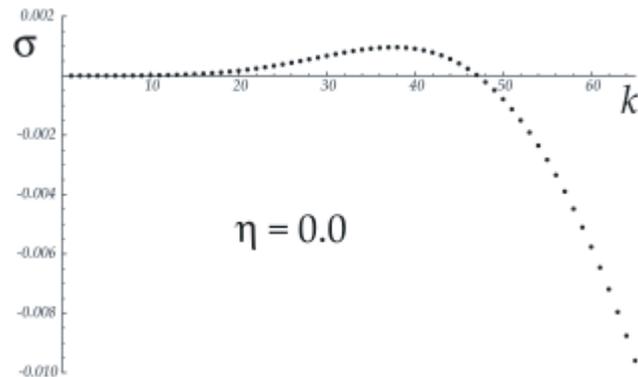
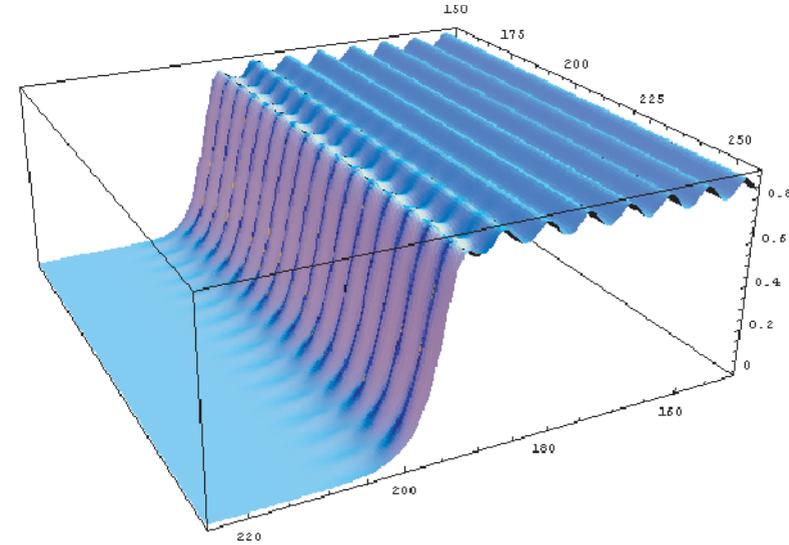
localized states in the interface

$$\partial_t P = \varepsilon P_{yy} + P_y^2 P_{yy} - P_{yyyy} + \alpha P_y^2 - \nu \sin(\kappa P)$$



Amended amplitud equation

- The spectrum of the interface



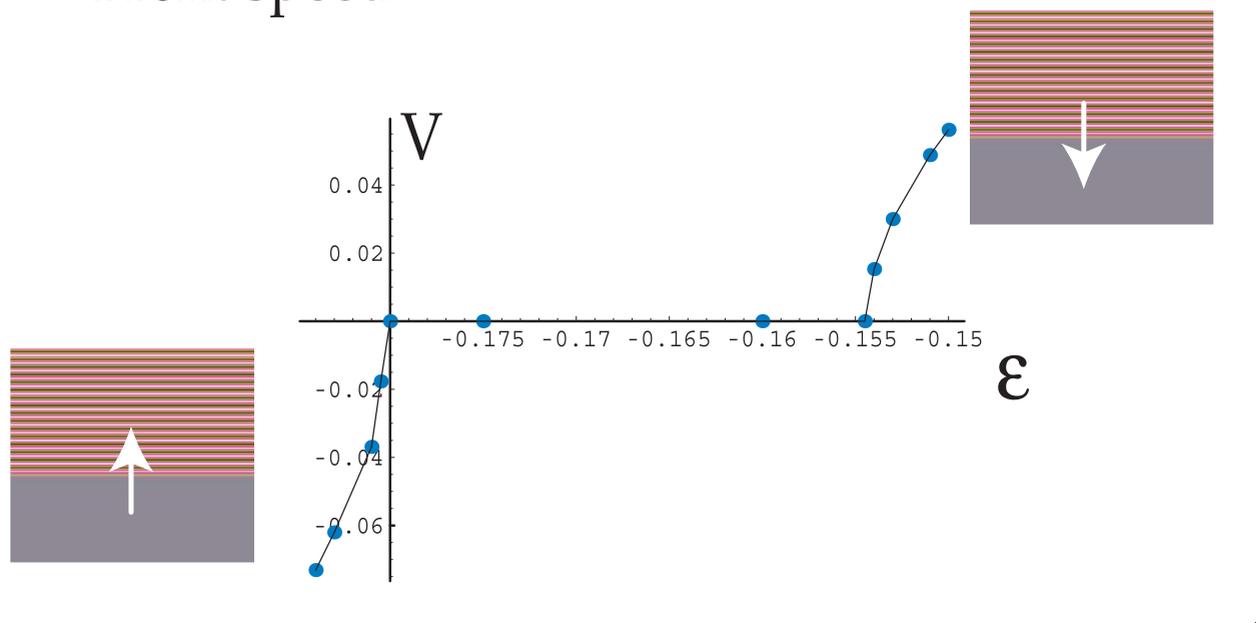
Hence, the needlework are consequence of the interaction of envelope variation with the small scale underlying the spatial periodic solution

Pinning effect in anisotropic system

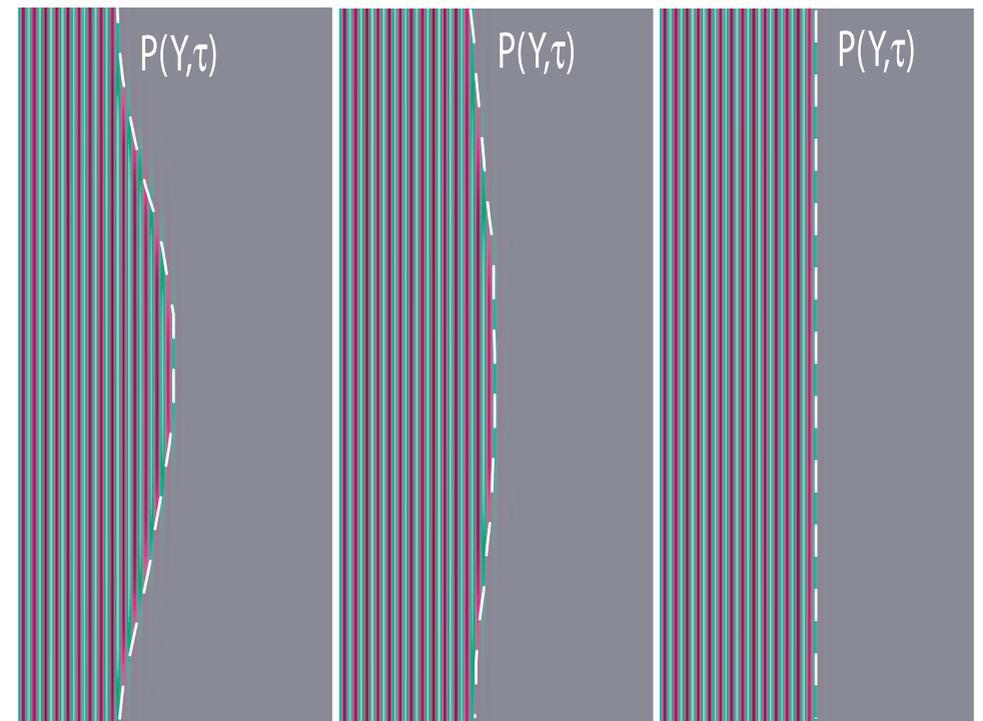
- The origin of the depinning in NewellWhiteheadSegel equation is the anisotropic spatial coupling. We consider the following model

$$\partial_t u = \varepsilon u + \nu u^3 - u^5 - (\partial_{xx} + q^2)^2 u + D \partial_{yy} u,$$

Front speed



Dynamical behavior

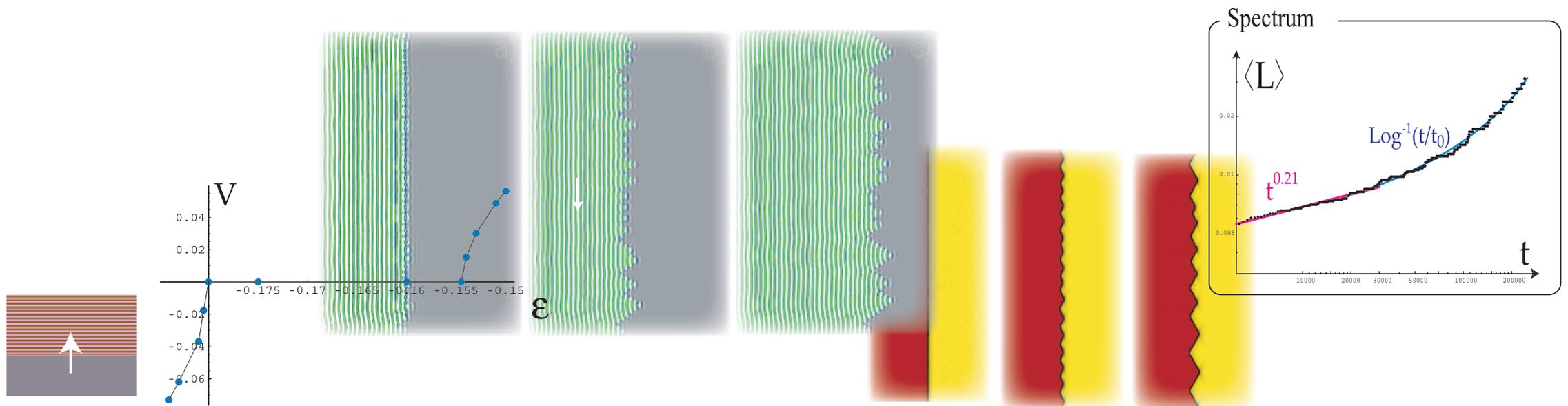


$$\partial_\tau B = \sigma B + |B|^2 B - |B|^4 B + \vec{\nabla}^2 B$$

Conclusions

Systems which have coexistence between stable stripe pattern and uniform states can exhibit interfaces connecting these states.

- The interface dynamics have complex behaviors: flat interface, periodic solutions, localized state, zigzag dynamics,....



Outlook

Noise induce front propagation

$$\partial_t u = \varepsilon u + \nu u^3 - u^5 - (\partial_{xx} + q^2)^2 u + D \partial_{yy} u, + \text{NOISE}$$

